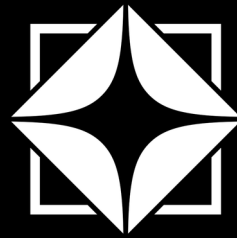


Heating/cooling of IGM due to Lyman- α Photons during Cosmic Dawn

Janakee Raste



NCRA • TIFR

with

Shiv Sethi (RRI) *and* Anjan Sarkar (NCRA-TIFR)

Advanced 21-cm Cosmology workshop
20th December, 2023

Outline

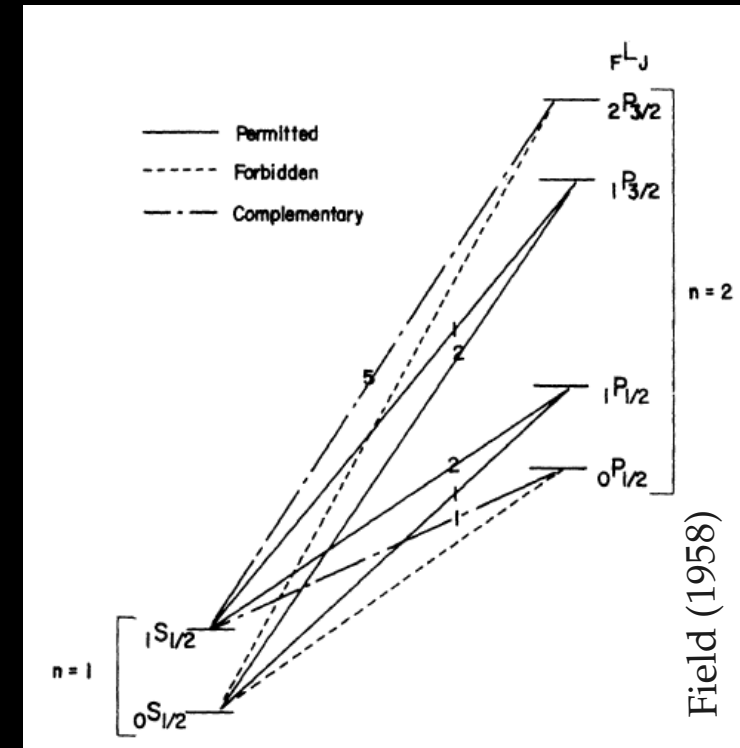
- The 21-cm Signal and Lyman- α photons cosmic dawn
- Injected and Continuum Lyman- α photons
- Shape of Lyman- α photons profiles and quasi-static solution
- Energy exchange between Lyman- α photons and HI atoms
- Coupled evolution of temperature and Lyman- α profile with time

Lyman- α Photons and 21-cm Signal from Cosmic Dawn

$$\Delta T_B \simeq 27 \text{ mK } x_{\text{HI}}(1 + \delta) \left(1 - \frac{T_{\text{CMB}}}{T_S}\right) \left(1 + \frac{1}{H} \frac{dv_p}{ds}\right)^{-1} \left(\frac{1+z}{10}\right)^{\frac{1}{2}} \left(\frac{0.14}{\Omega_m h^2}\right)^{\frac{1}{2}} \left(\frac{\Omega_b h^2}{0.022}\right)$$

- Photons with frequencies between Lyman- α and Lyman-limit escape the source halo and free stream into the medium
- When they redshifts into one of the Lyman series lines, they are absorbed by the HI atoms in the IGM.
- They undergo multiple scatterings before being redshifted out of the resonance width. These scatterings couple the spin temperature T_S to T_K :

$$T_S = \frac{T_{\text{CMB}} + y_c T_K + y_\alpha T_\alpha}{1 + y_c + y_\alpha}$$



$$y_\alpha \simeq 5.9 \times 10^{11} \frac{n_\alpha}{T_K^{3/2}}$$

Time evolution of Lyman- α photons intensity profile

- Diffusion equation with Fokker-Planck approximation:

$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \gamma \frac{\partial J}{\partial x} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

Time evolution of Lyman- α photons intensity profile

- Diffusion equation with Fokker-Planck approximation:

$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \gamma \frac{\partial J}{\partial x} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

diffusion term

redshift

recoil

source function

Time evolution of Lyman- α photons intensity profile

- Diffusion equation with Fokker-Planck approximation:

$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \gamma \frac{\partial J}{\partial x} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

diffusion term

redshift

recoil

source function

$$x = \frac{\nu - \nu_\alpha}{\Delta\nu_D}$$

$$\Delta\nu_D = \nu_\alpha \frac{v_T}{c} = \sqrt{\frac{2k_B T_K}{m_p c^2}} \nu_\alpha$$

$$\eta = \frac{h_p \nu_\alpha}{c \sqrt{2k_B T_K m_p}}$$

$$\phi(x, a) = \frac{a}{\pi^{3/2}} \int_{-\infty}^{\infty} \frac{\exp(-y^2)}{a^2 + (x - y)^2} dy$$

$$\gamma = \frac{1}{\tau_{GP}} = \frac{m_e \nu_\alpha}{\pi e^2 f_\alpha} \frac{H(z)}{n_{HI}}$$

$$a = \frac{A_\alpha}{4\pi \Delta\nu_D}$$

Time evolution of Lyman- α photons intensity profile

- Diffusion equation with Fokker-Planck approximation:

$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \gamma \frac{\partial J}{\partial x} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

diffusion term

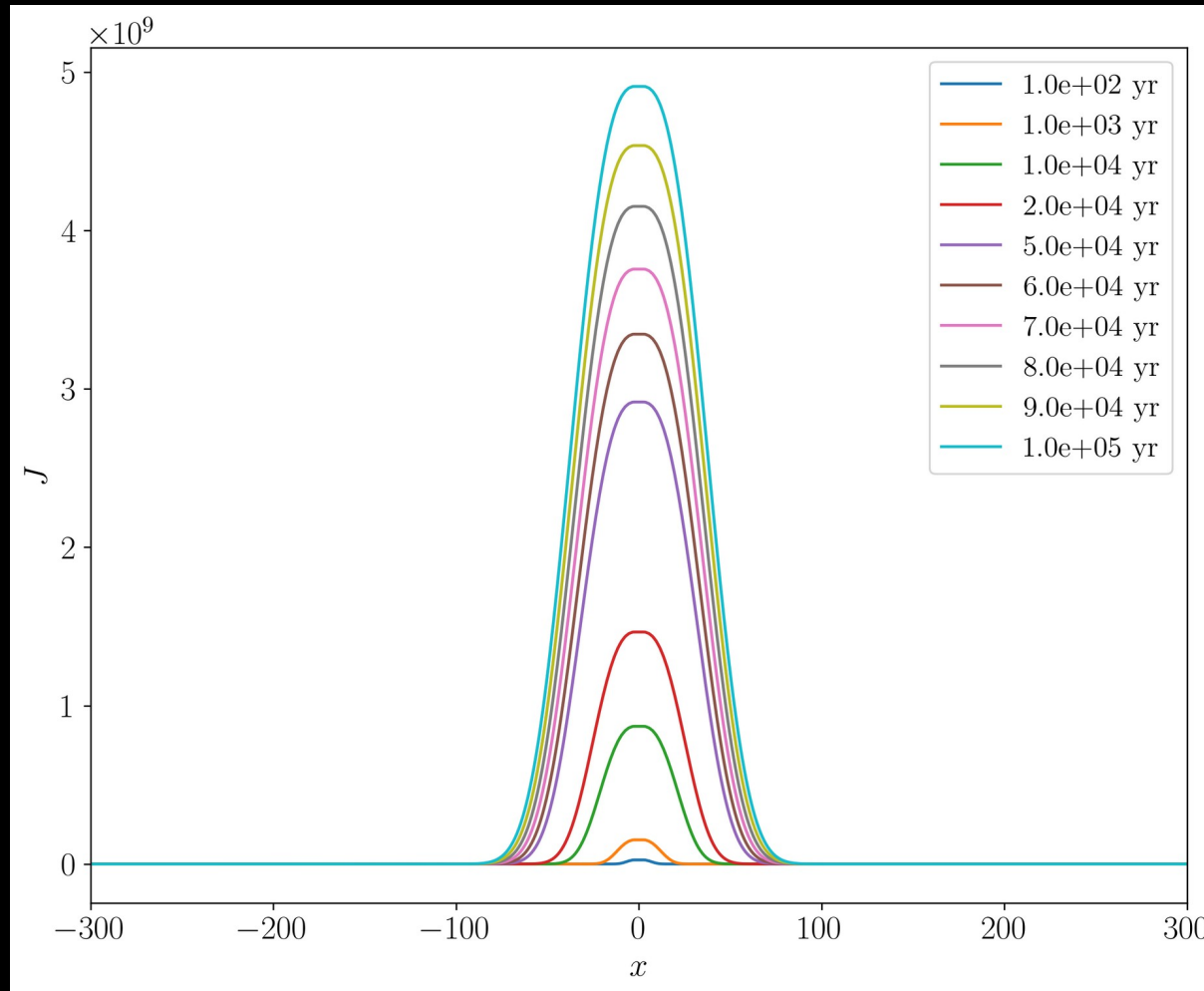
redshift

recoil

source function

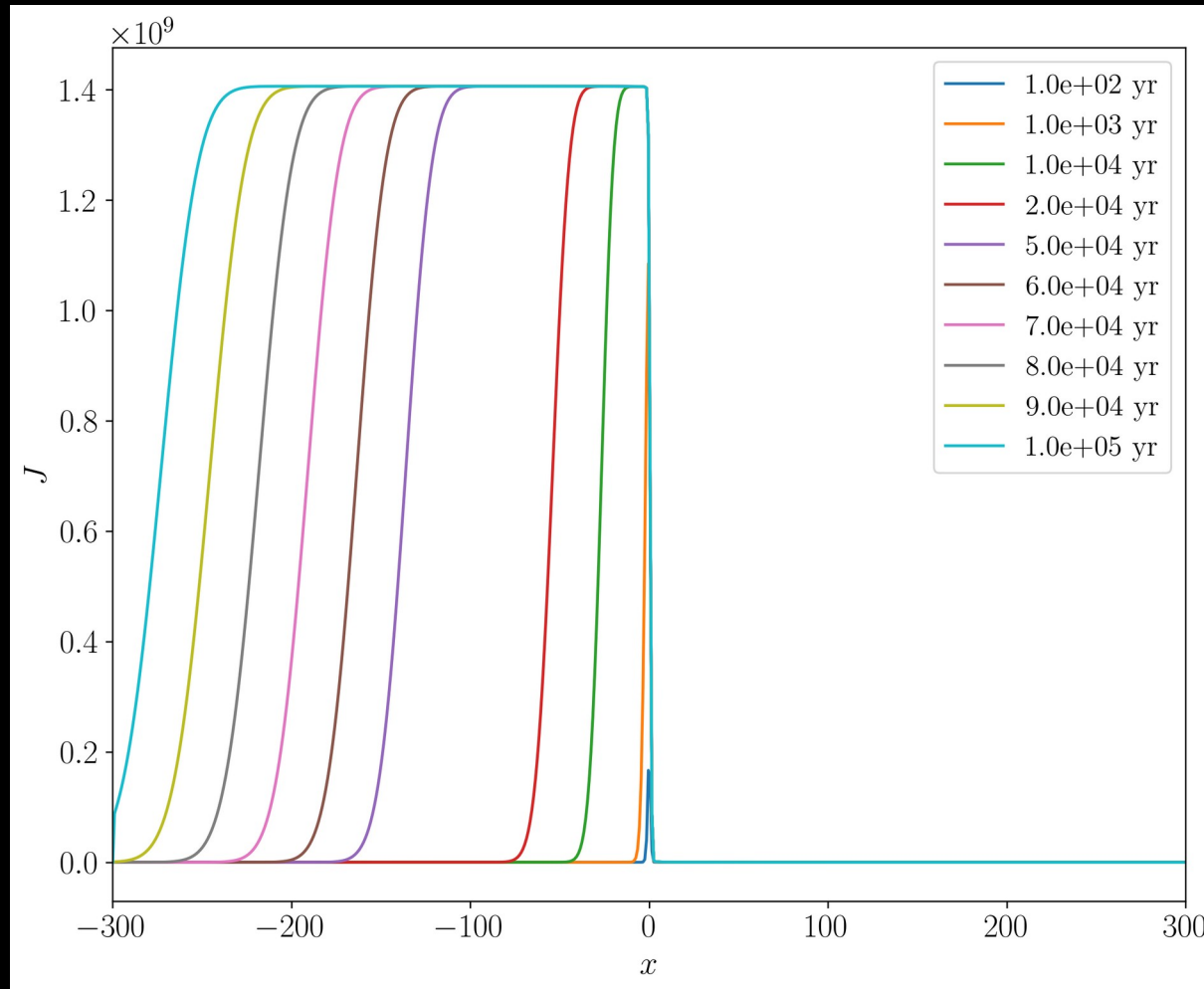
- **Continuum photons:**
 - Photons which have redshifted into the Lyman- α line
- **Injected photons:**
 - Photons which are injected at the Lyman- α line centre due to de-excitation of HI atoms

Injected Photons profile shape



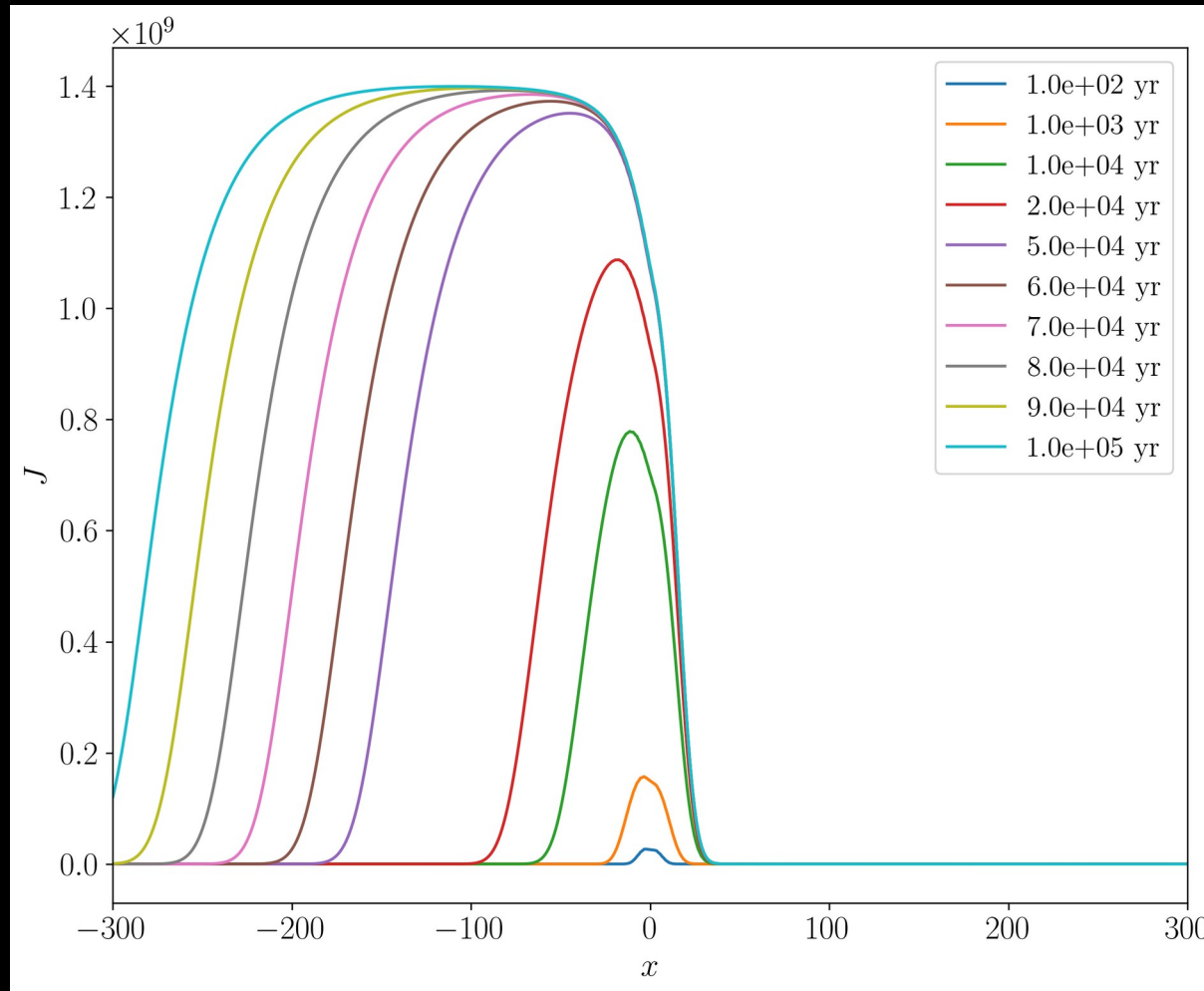
$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \cancel{\gamma \frac{\partial J}{\partial x}} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

Injected Photons profile shape



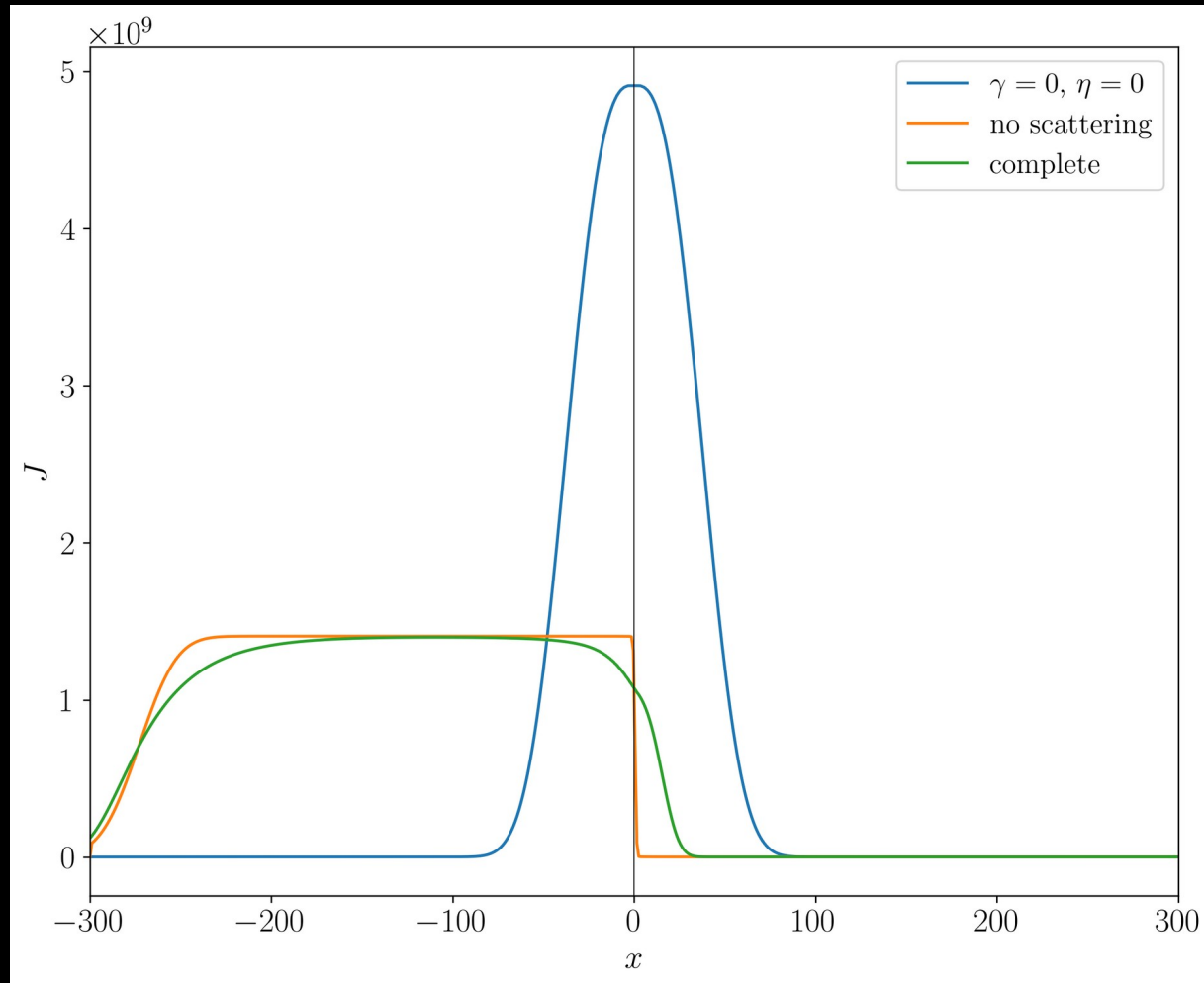
$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \gamma \frac{\partial J}{\partial x} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

Injected Photons profile shape

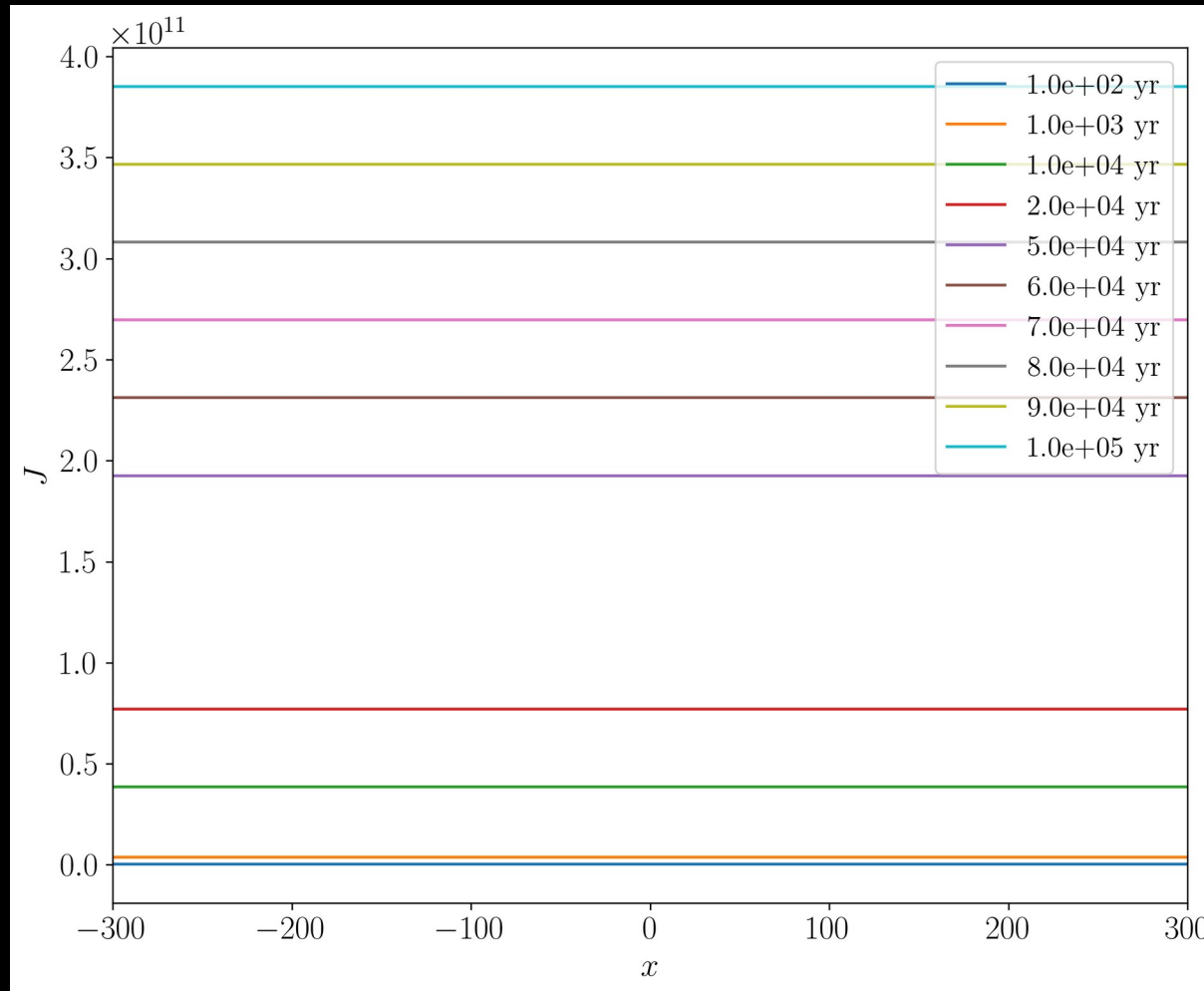


$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \gamma \frac{\partial J}{\partial x} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

Injected Photons profile shape

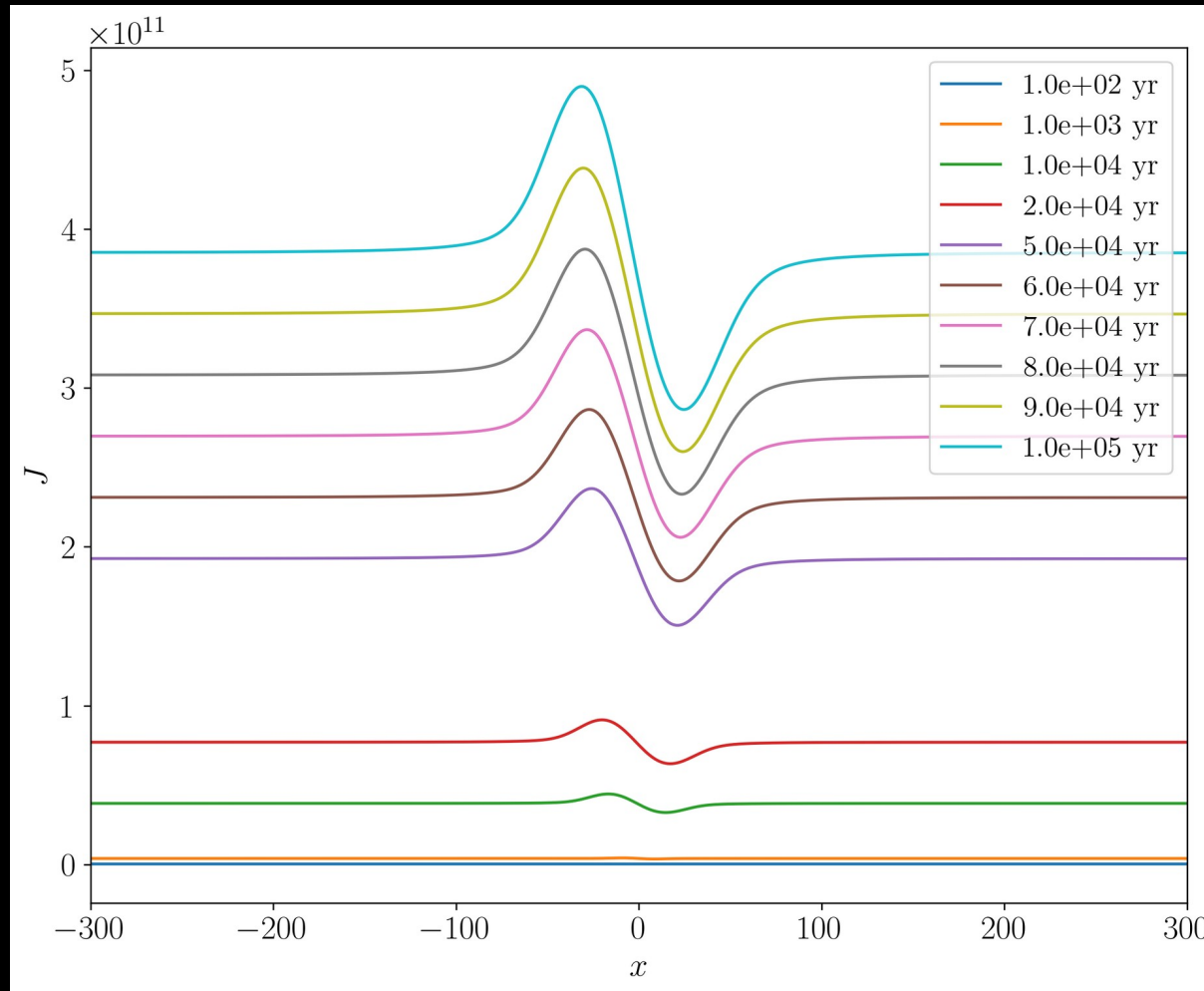


Continuum Photons profile shape



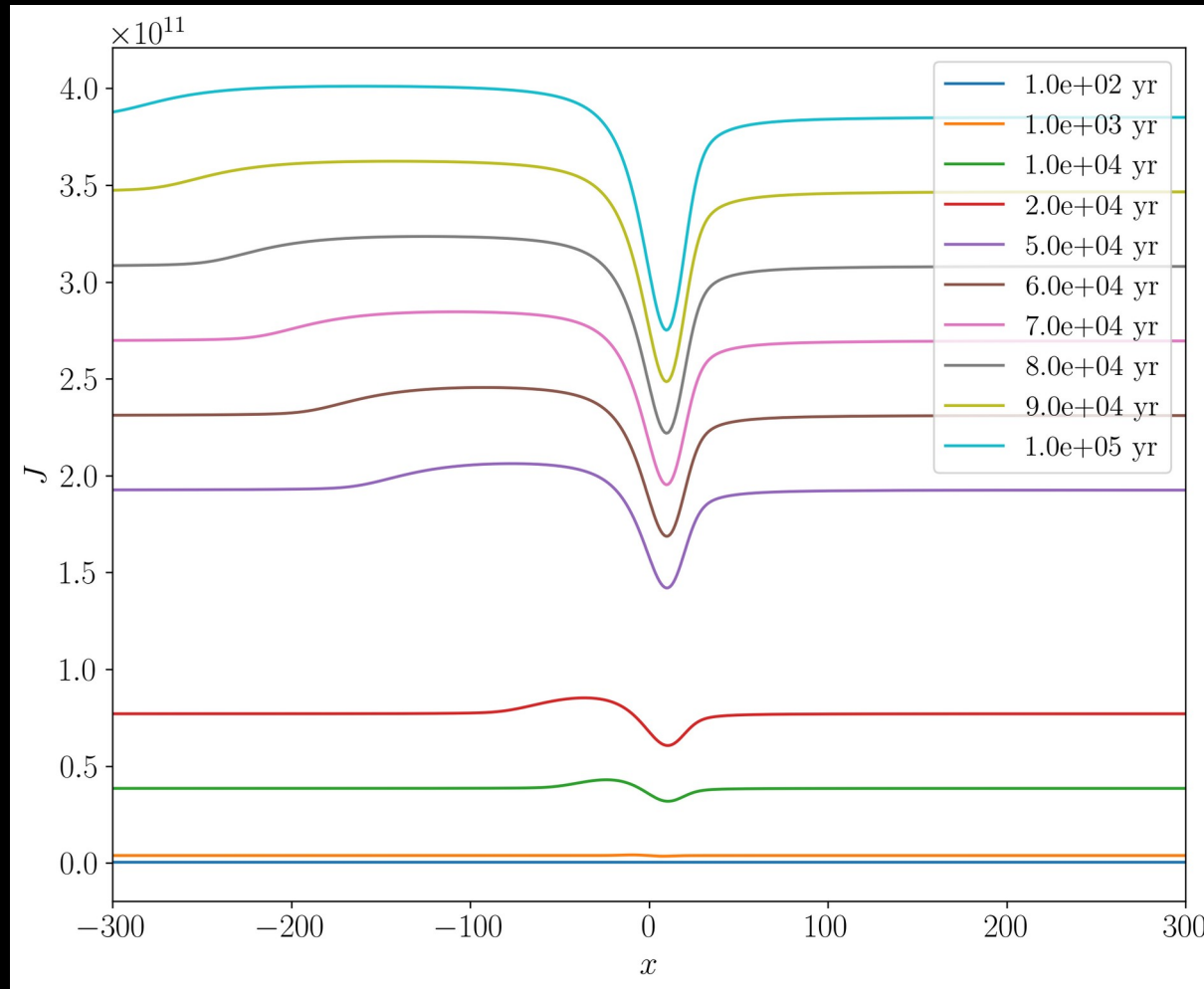
$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \gamma \frac{\partial J}{\partial x} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

Continuum Photons profile shape



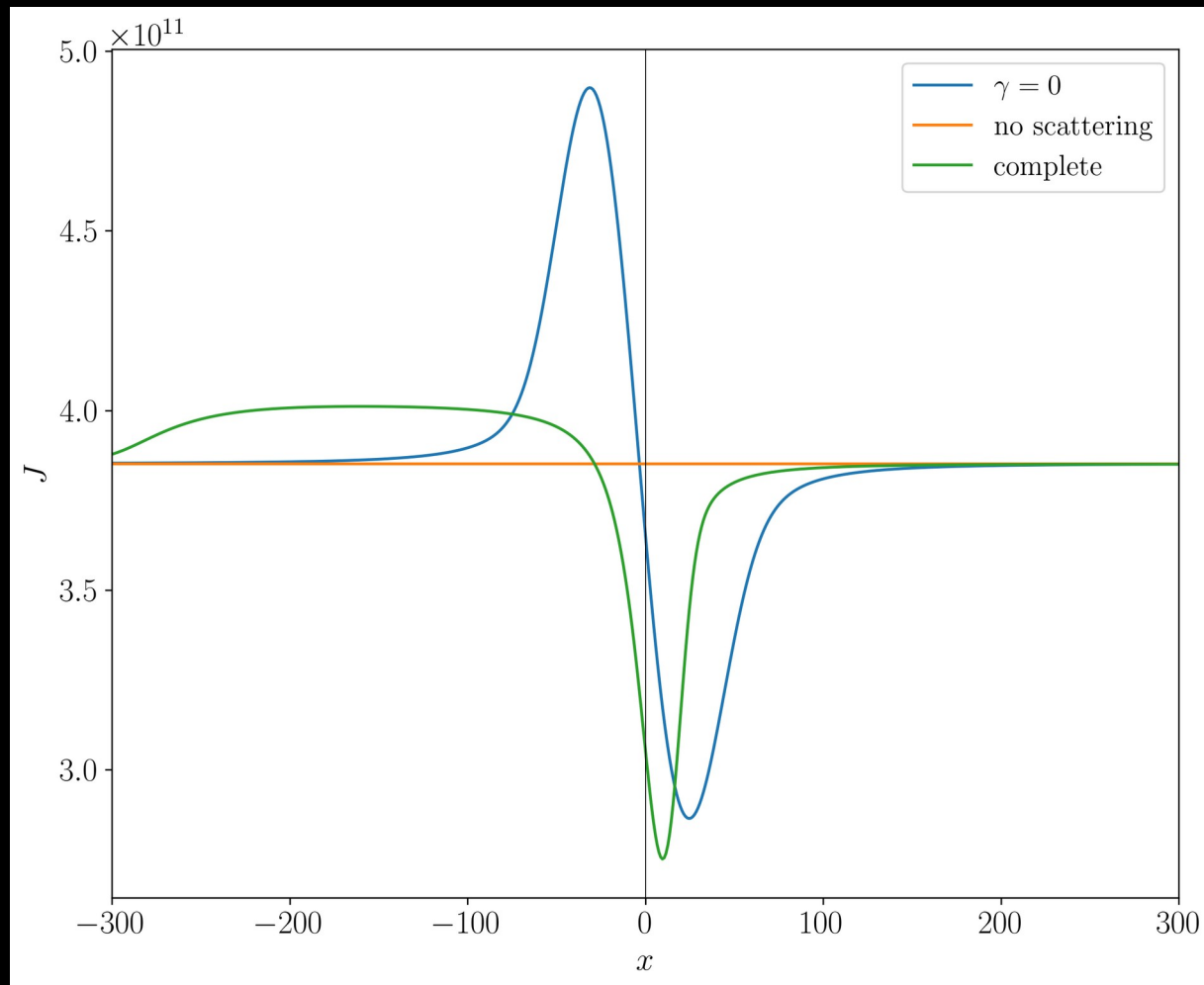
$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \cancel{\gamma \frac{\partial J}{\partial x}} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

Continuum Photons profile shape



$$\frac{\partial J(x, t)}{\partial t} = \frac{1}{2} \frac{\partial}{\partial x} \left[\phi(x) \frac{\partial J}{\partial x} \right] + \gamma \frac{\partial J}{\partial x} + \eta \frac{\partial}{\partial x} [\phi J] + C(t)$$

Continuum Photons profile shape



Energy exchange between Lyman- α Photons and HI atoms

- During the repeated scattering process, the Lyman- α photons also exchange energy with the medium.
 - Compton and inverse-compton scattering between Lyman- α photons and HI atoms
 - Wouthuysen-Field effect in which the energy exchange occurs owing to the change in the level populations of hyperfine states

$$\dot{Q} = \int \frac{(h\nu_\alpha)^2}{m_p c^2} \sigma_0 \phi \left(J + \frac{kT_k}{\Delta\nu_D h} J' \right) dx + \int \frac{(h\nu_{21})^2}{2kT_s} \sigma_0 \phi \left(J + \frac{kT_s}{\Delta\nu_D h} J' \right) dx$$

- **Continuum photons** usually heat up the medium
- **Injected photons** usually cool the medium

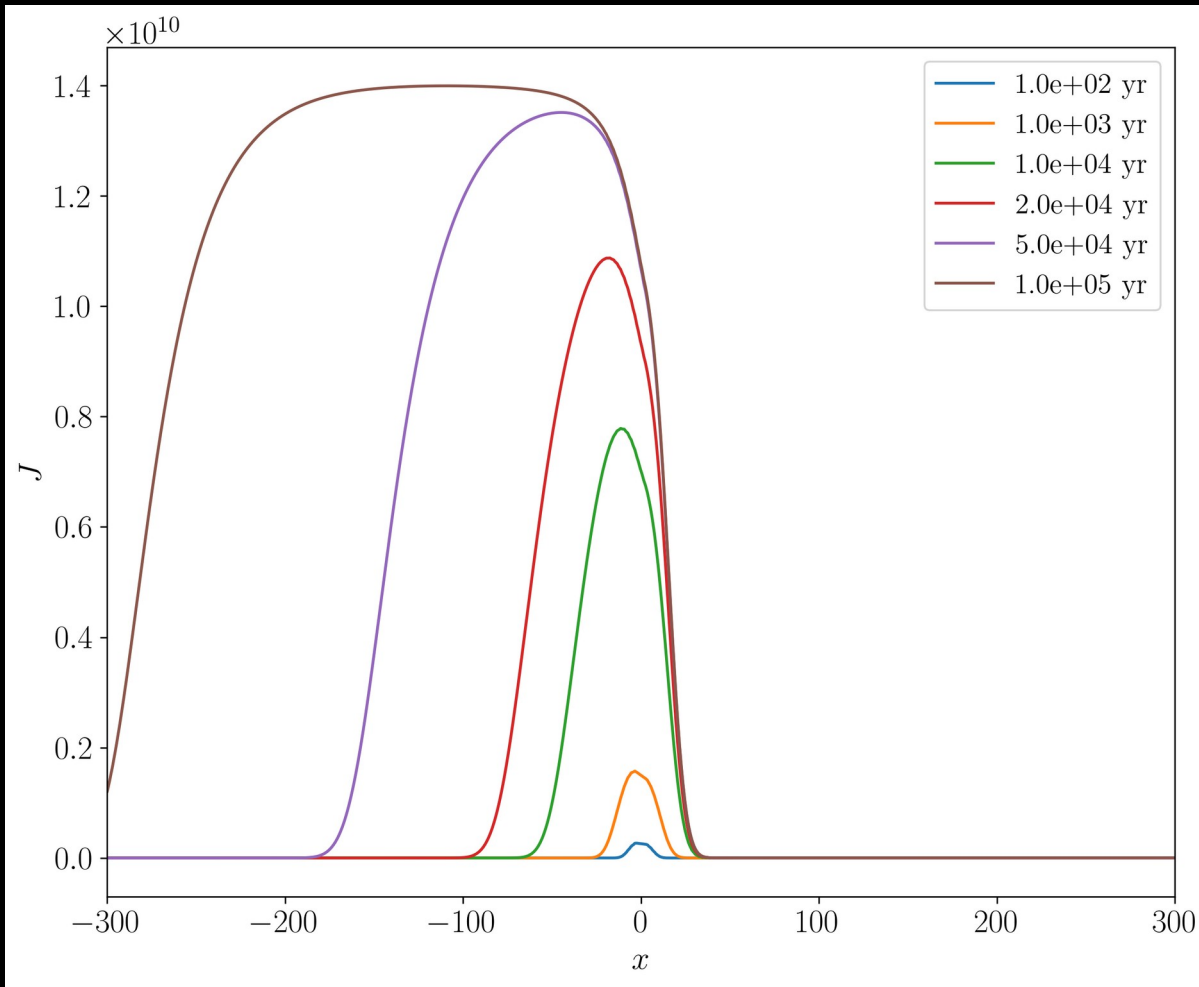
Energy exchange between Lyman- α Photons and HI atoms

- During the repeated scattering process, the Lyman- α photons also exchange energy with the medium.
 - Compton and inverse-compton scattering between Lyman- α photons and HI atoms
 - Wouthuysen-Field effect in which the energy exchange occurs owing to the change in the level populations of hyperfine states

$$\dot{Q} = \int \frac{(h\nu_\alpha)^2}{m_p c^2} \sigma_0 \phi \left(J + \frac{kT_k}{\Delta\nu_D h} J' \right) dx + \int \frac{(h\nu_{21})^2}{2kT_s} \sigma_0 \phi \left(J + \frac{kT_s}{\Delta\nu_D h} J' \right) dx$$

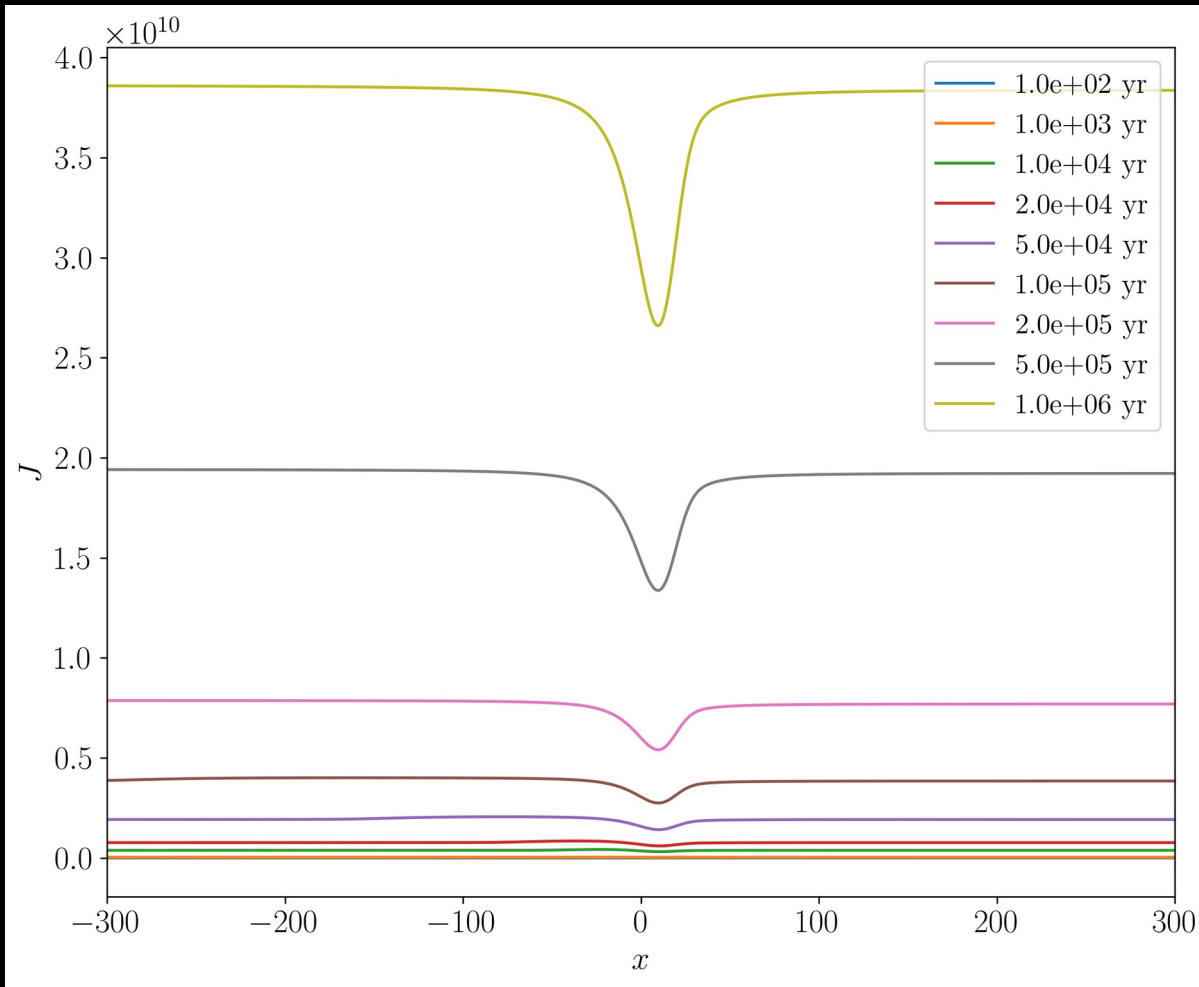
- **Continuum photons** usually heat up the medium
- **Injected photons** usually cool the medium
- **Caveat:** Quasi-static equilibrium of the Lyman- α profile is assumed

Why do Injected Photons Cool?



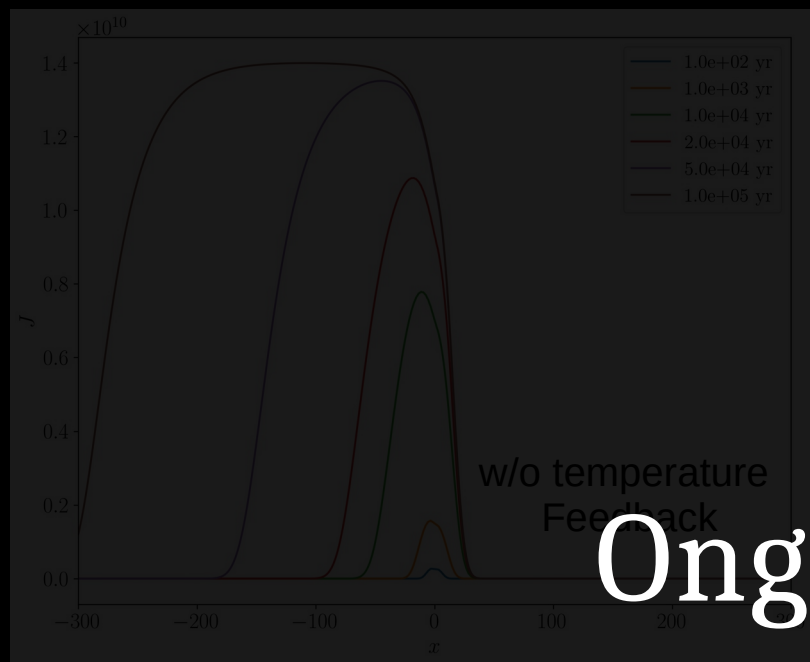
- Photons injected with Voigt profile at the Lyman- α line center.
- Due to redshifting, more redward photons than blueward photons
- More red \rightarrow blue scattering
- More cooling

Why do Continuum Photons Heat?

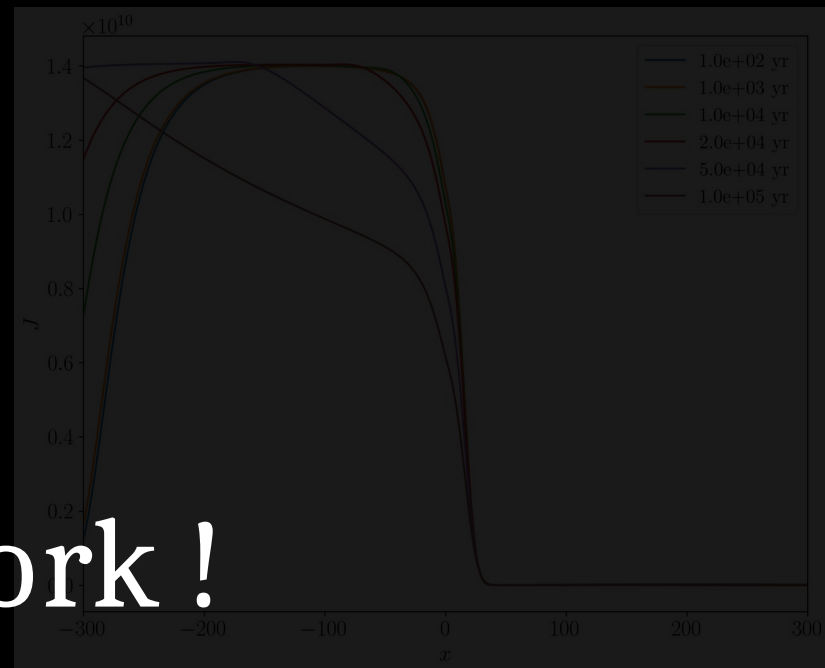


- Photons redshifting into the Lyman- α line.
- More new blueward photons due to redshift
- More blue \rightarrow red scattering
- More heating

Fiducial Injected Ly- α profile evolution for $z = 20$ and $T_K = 10$ K

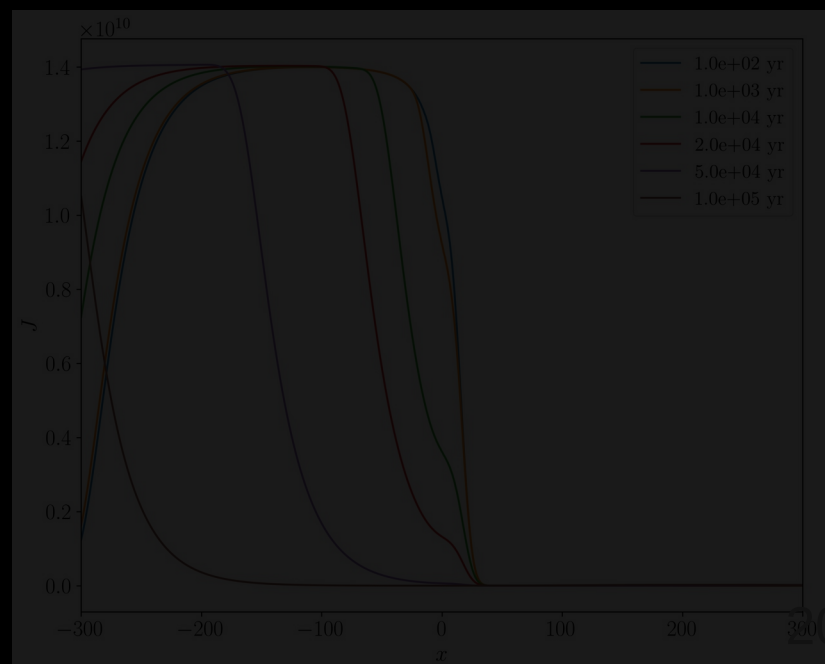
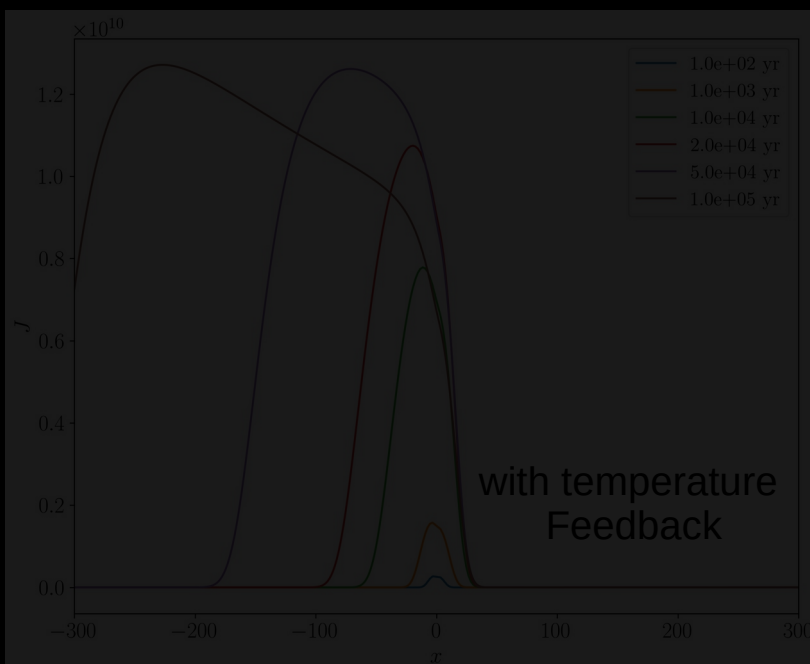


Temperature Feedback starts

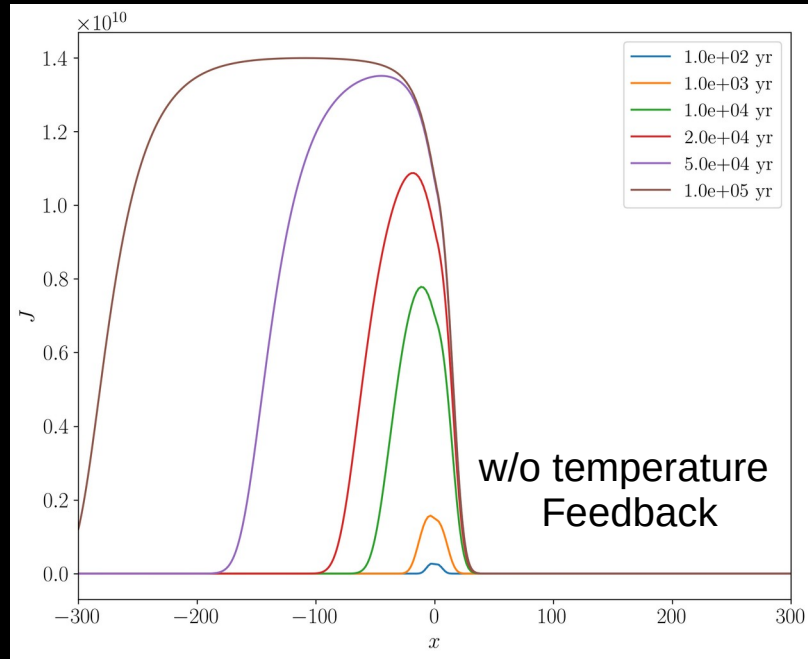


Ongoing work !

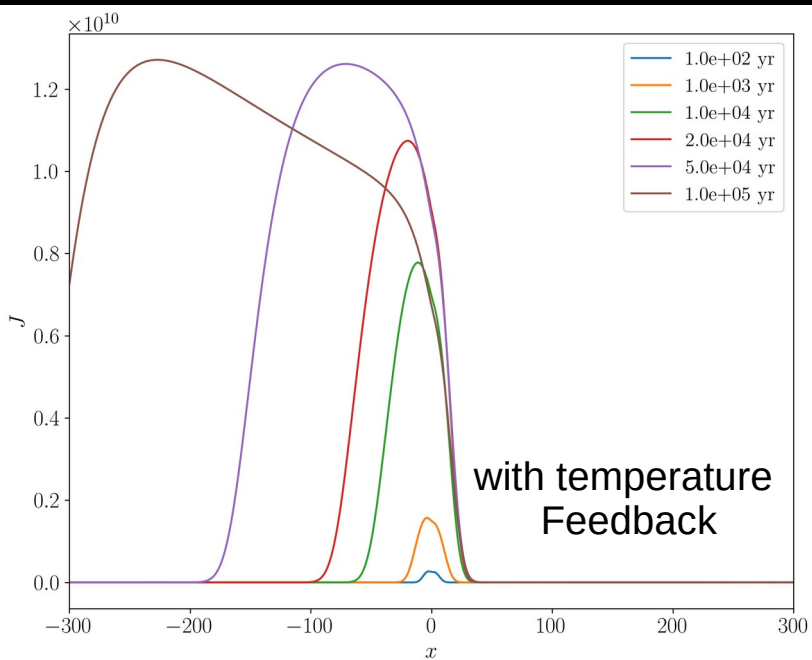
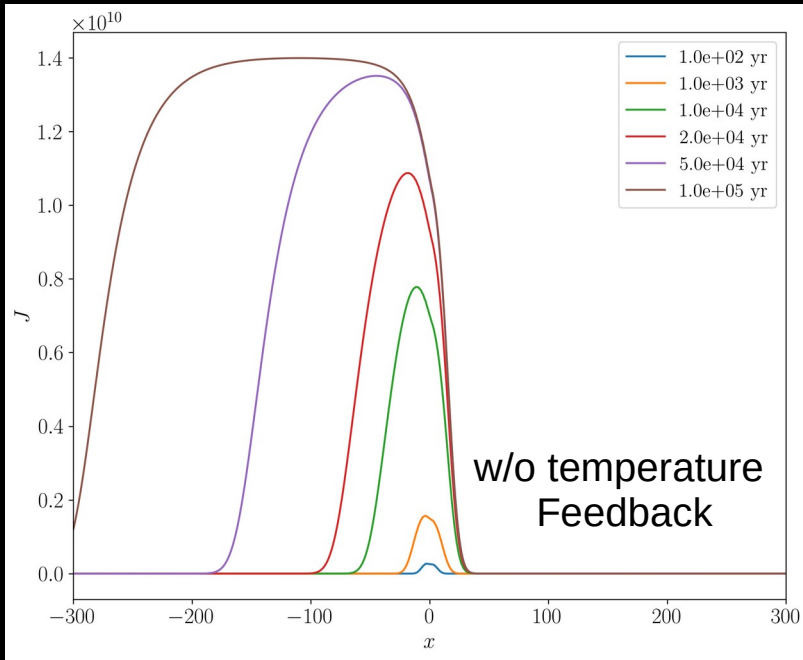
Source switches off



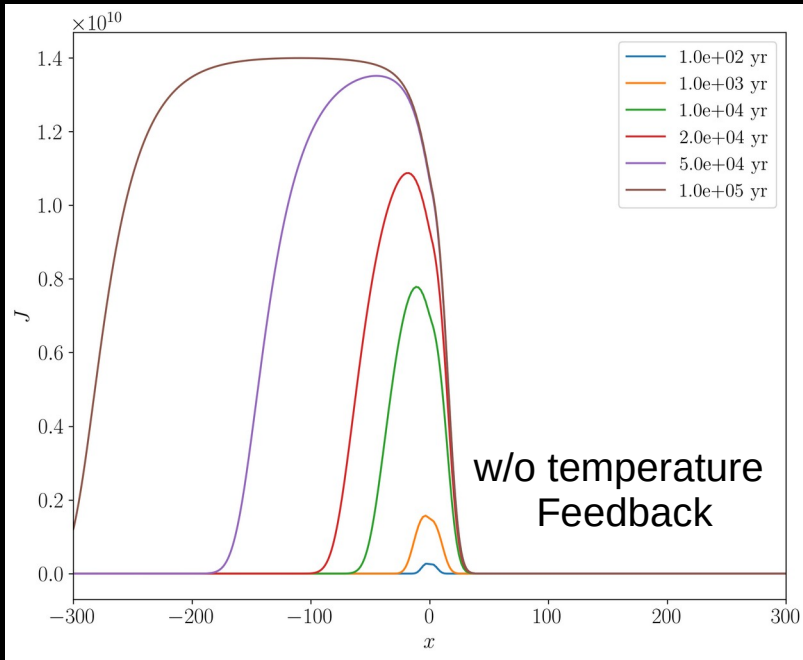
Fiducial **Injected** Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



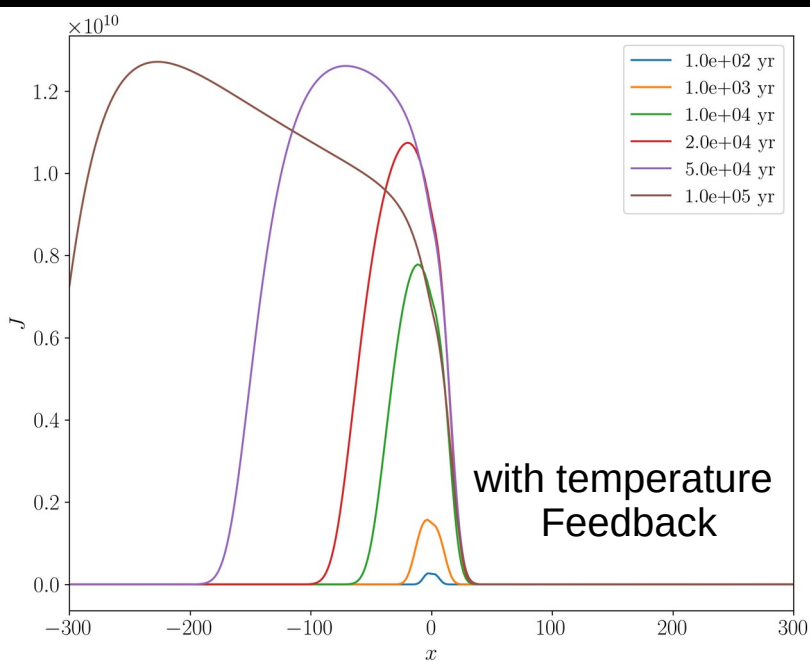
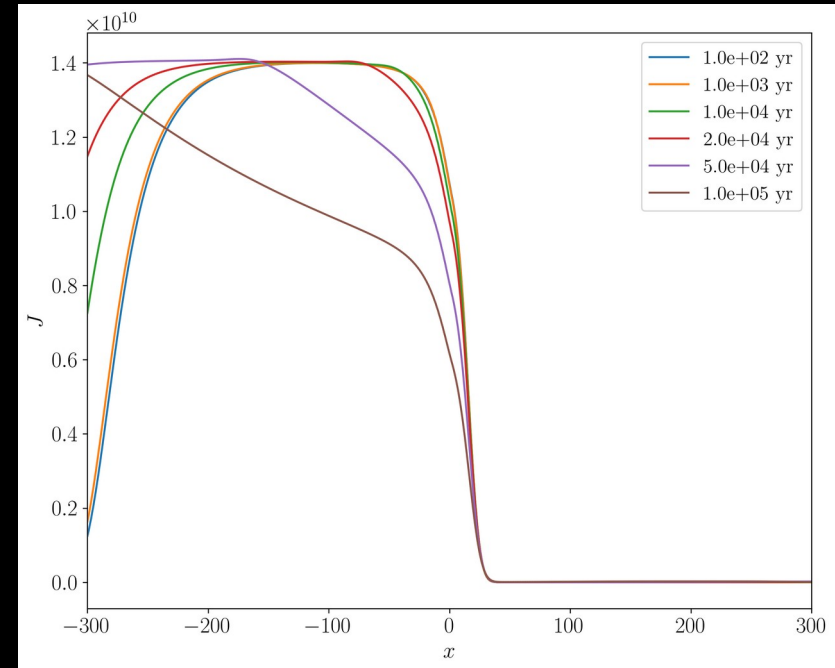
Fiducial **Injected** Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



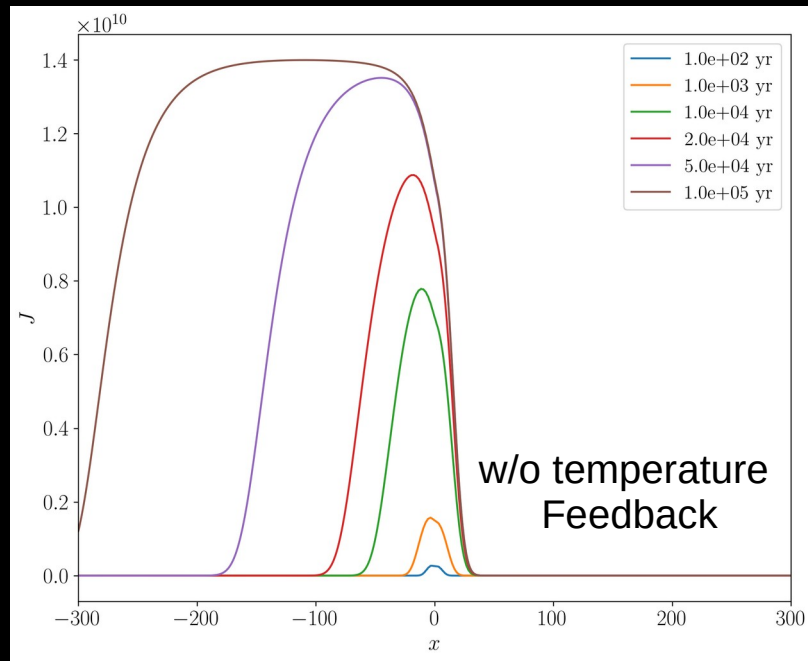
Fiducial **Injected** Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



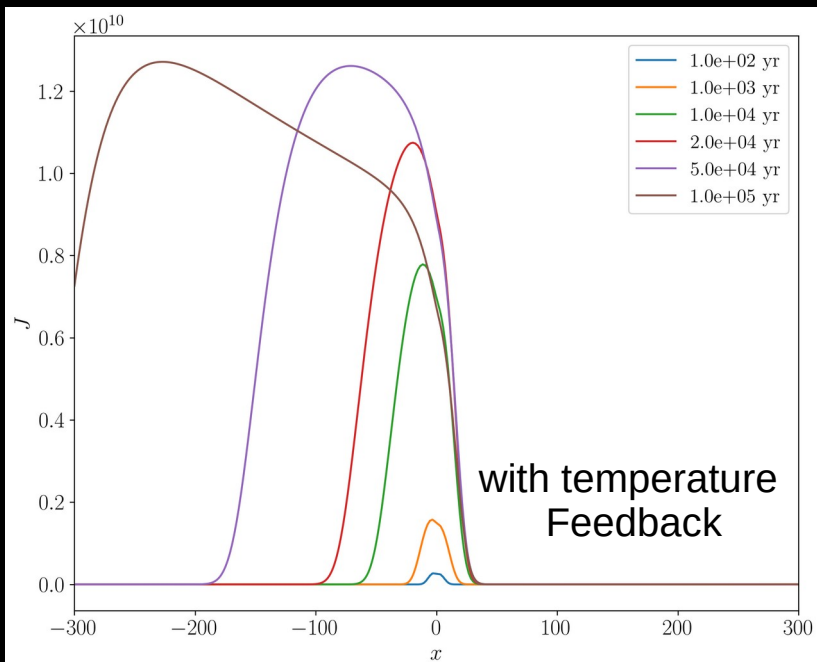
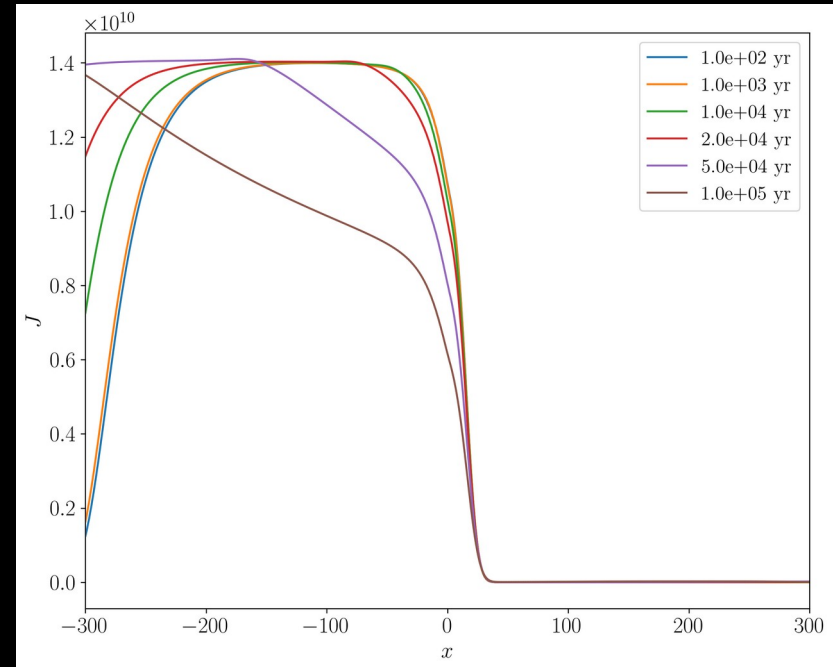
Temperature Feedback starts



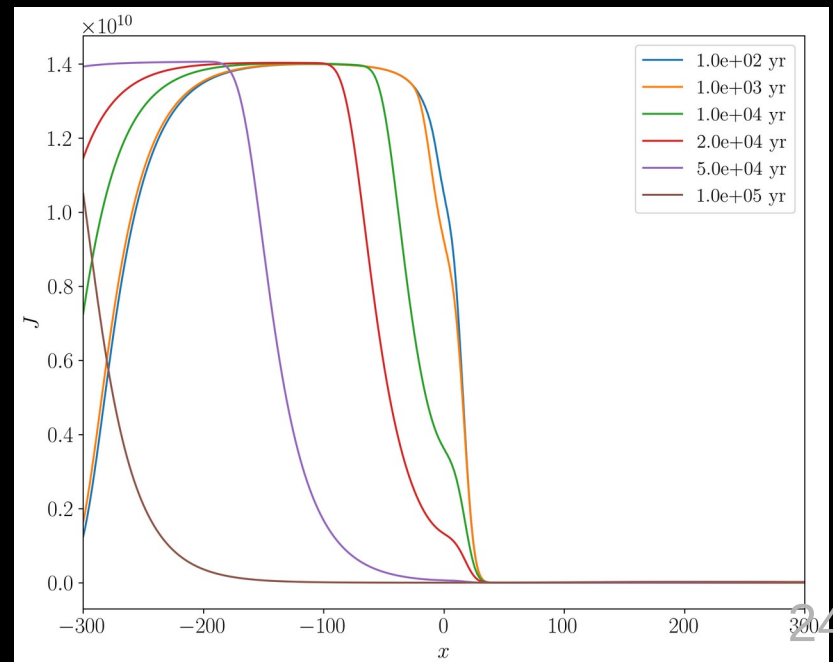
Fiducial **Injected** Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



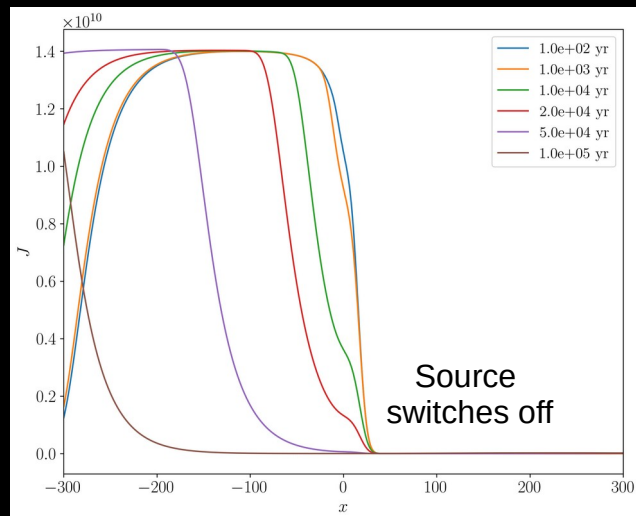
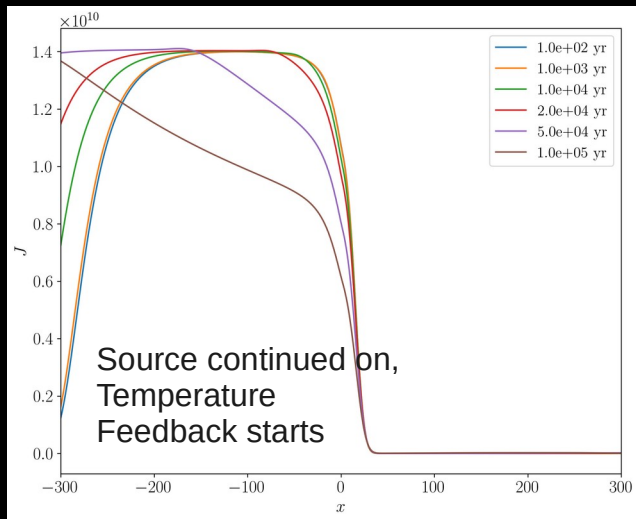
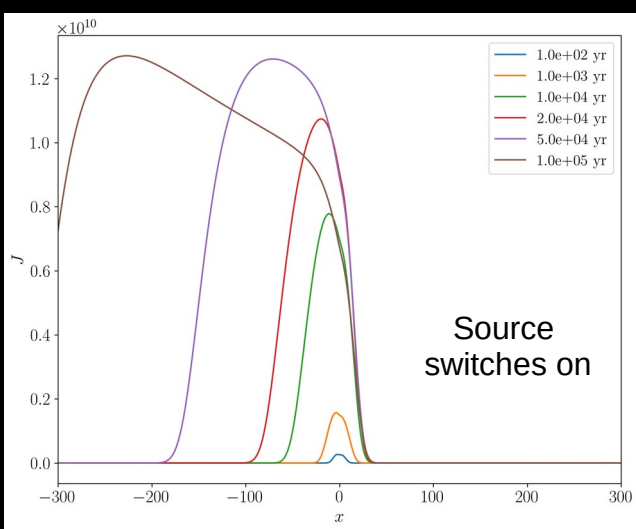
Temperature Feedback starts



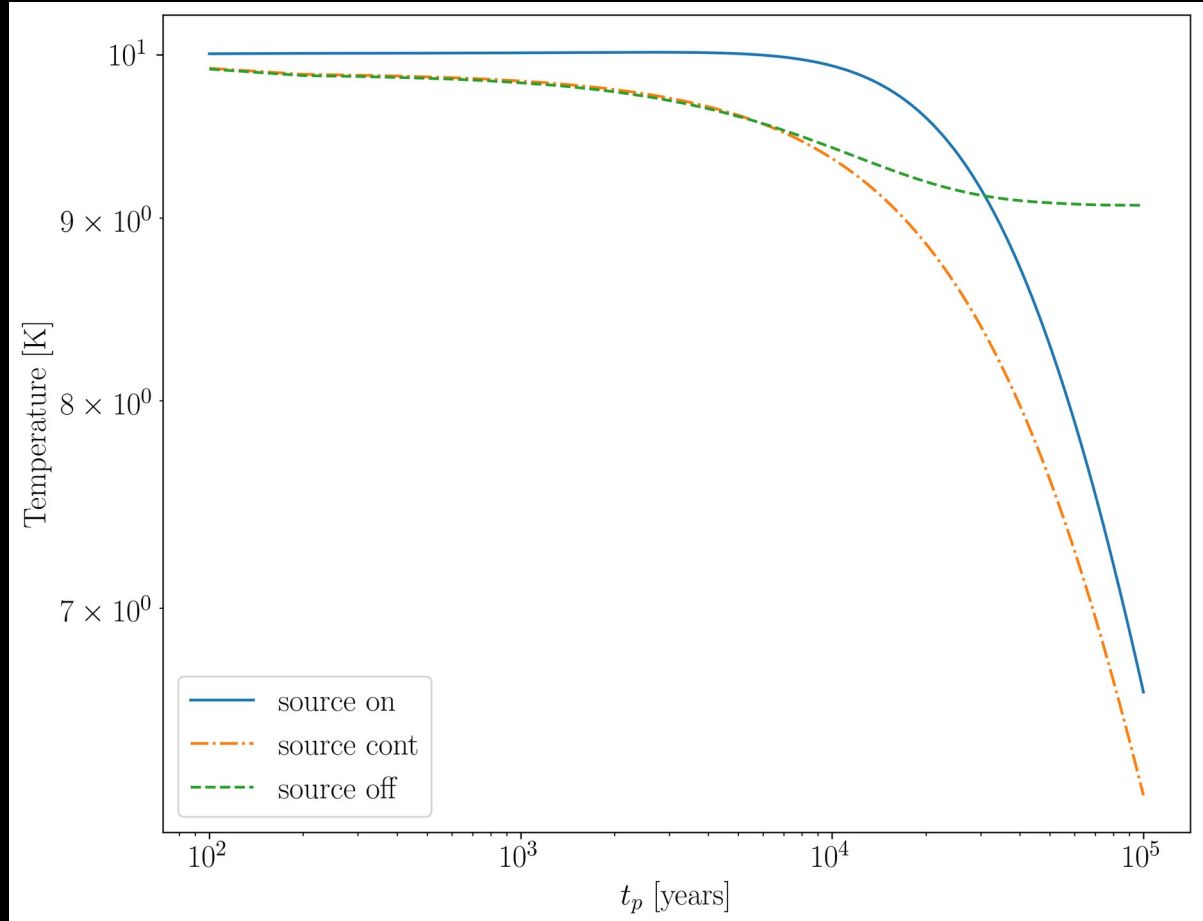
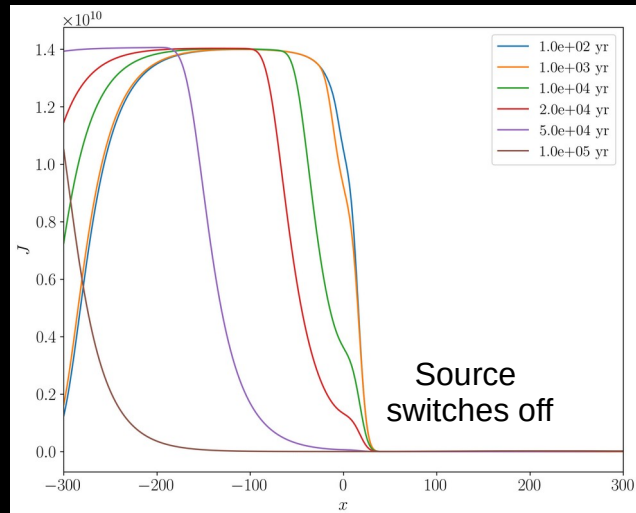
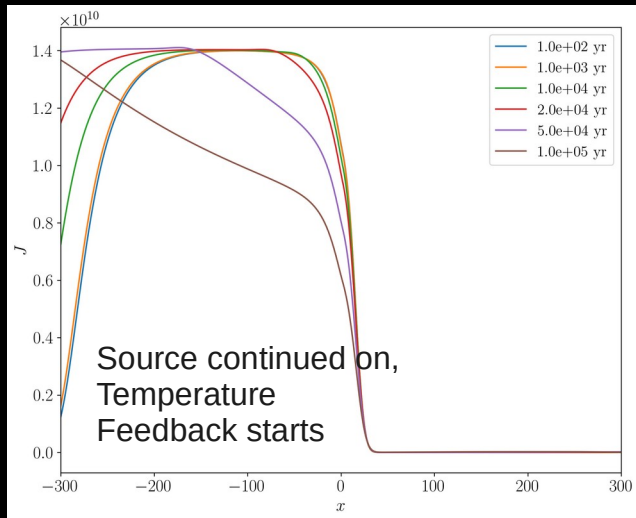
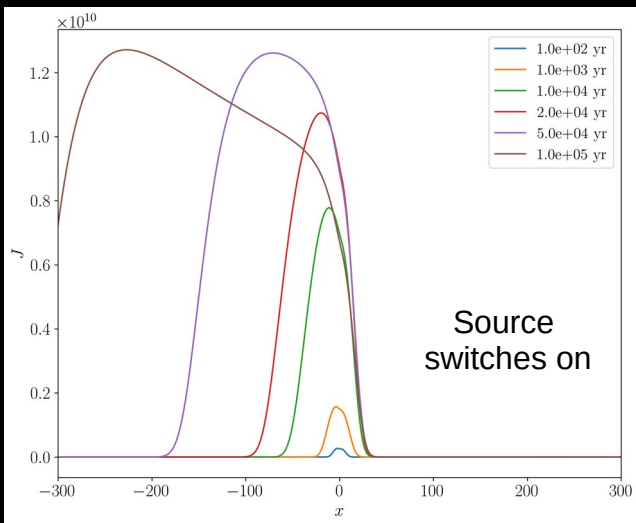
Source switches off



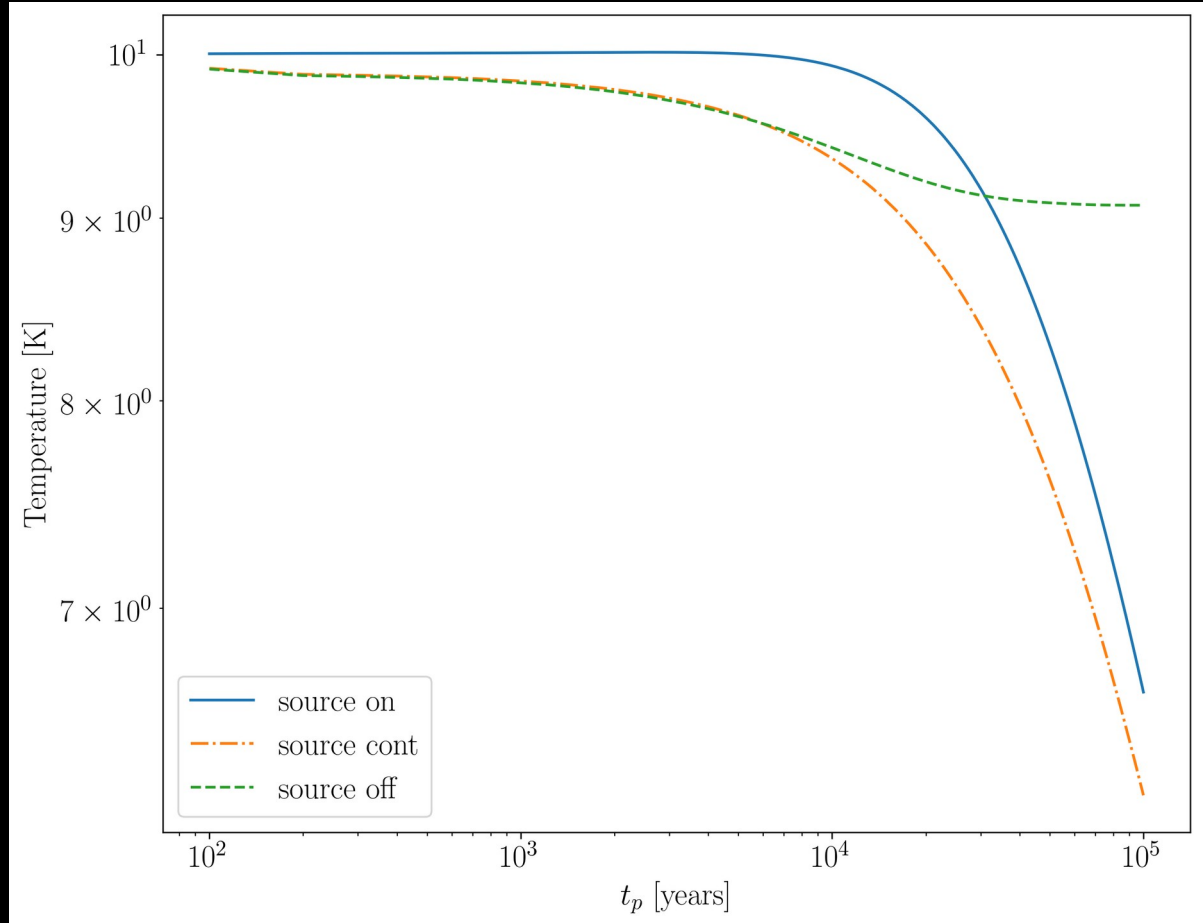
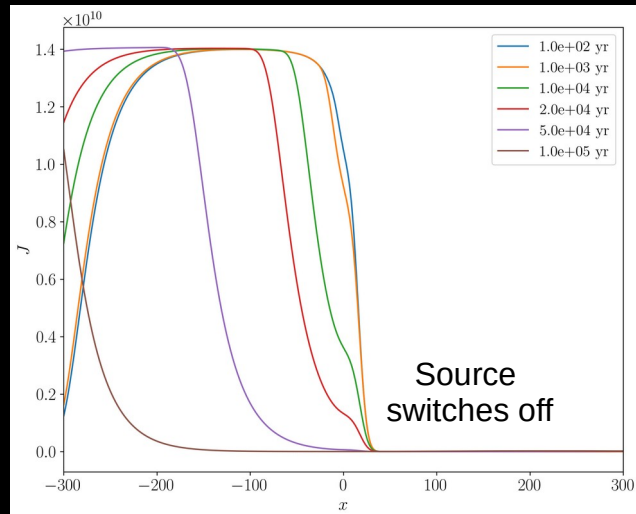
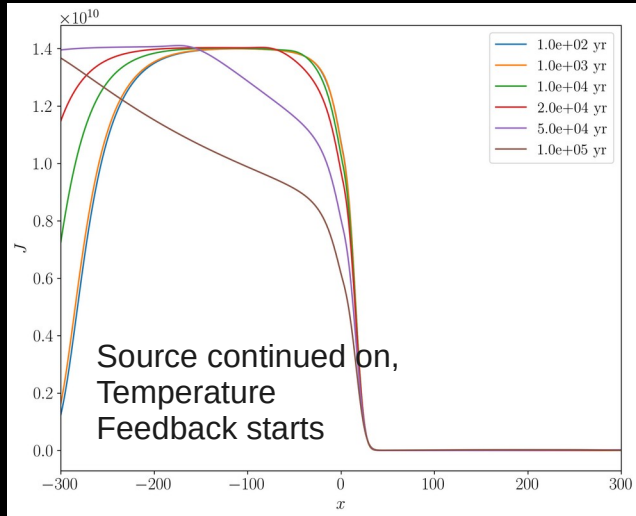
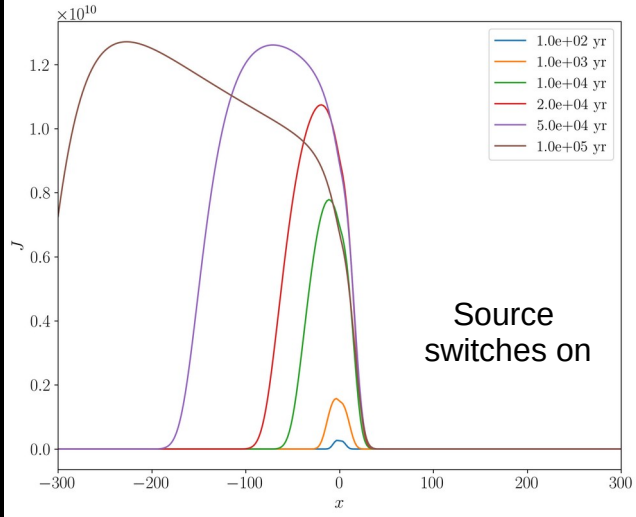
Fiducial Injected Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



Fiducial Injected Ly- α profile evolution for $z = 20$ and $T_K = 10$ K

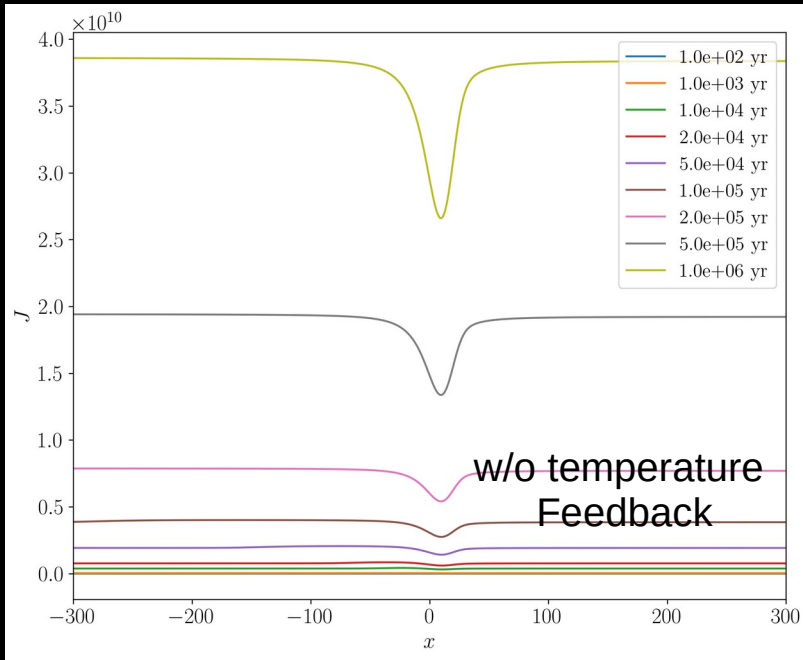


Fiducial Injected Ly- α profile evolution for $z = 20$ and $T_K = 10$ K

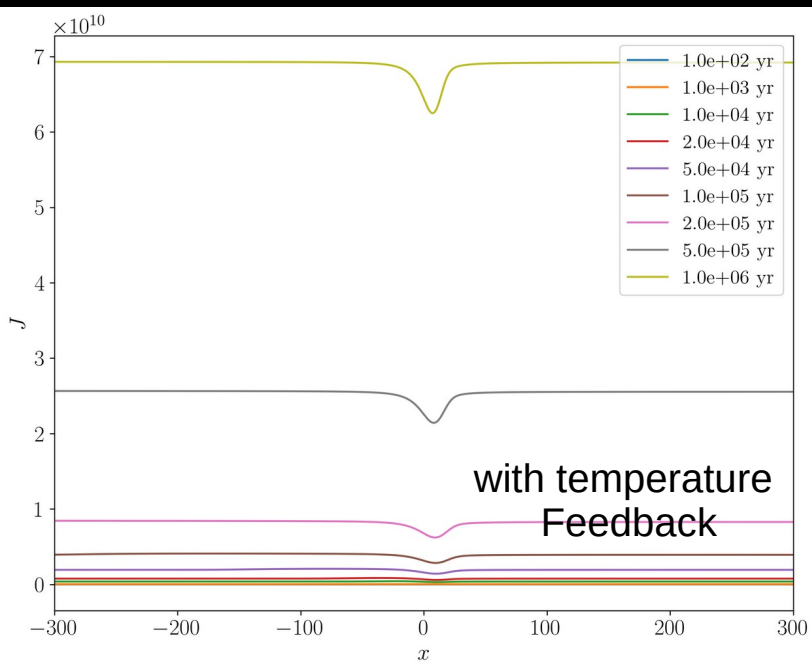
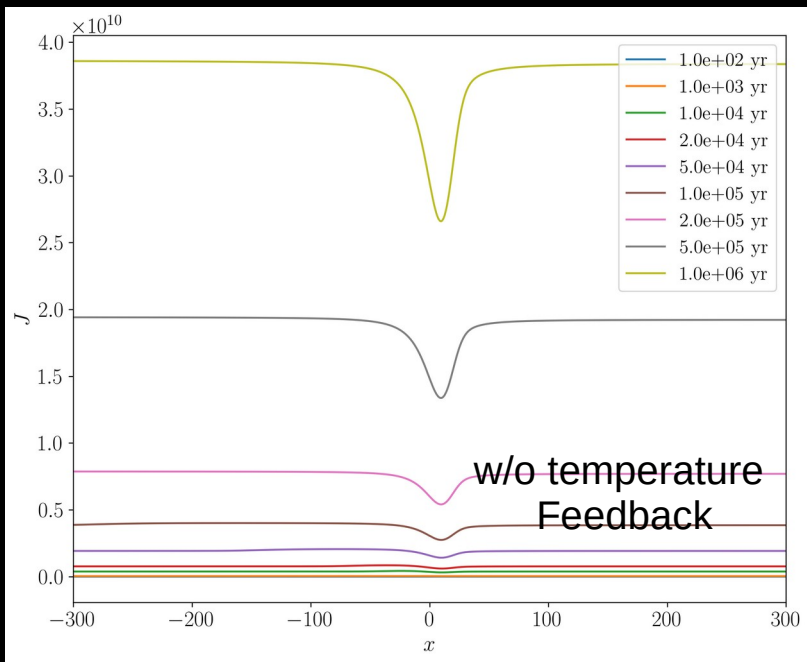


- Temperature feedback, source switching on, switching off delay/prevent quasi-static solution
- Injected photons cause overall less cooling than previously computed in literature.

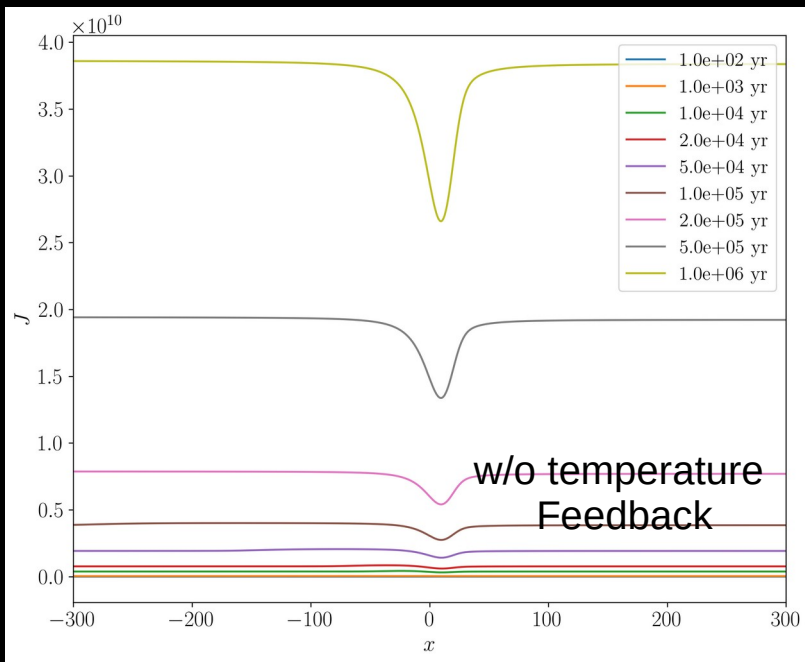
Fiducial Continuum Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



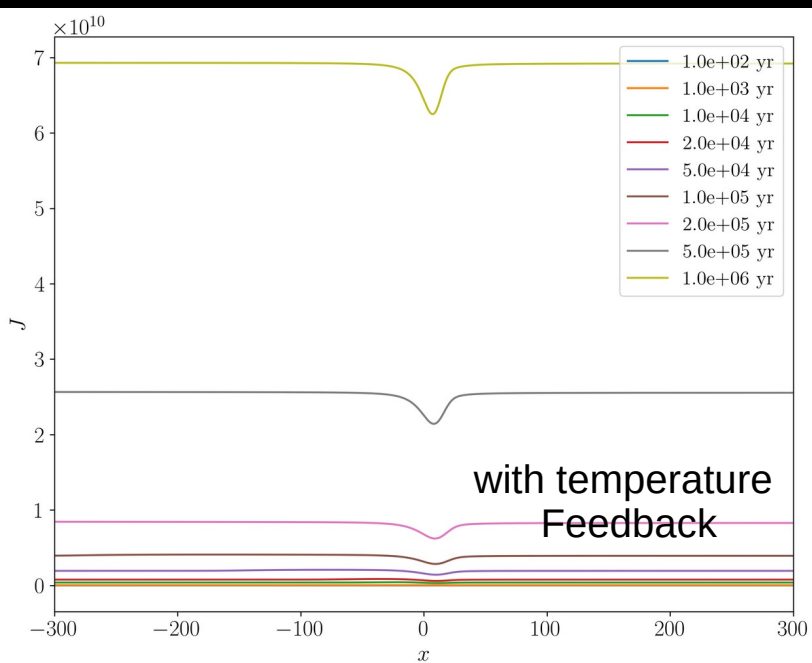
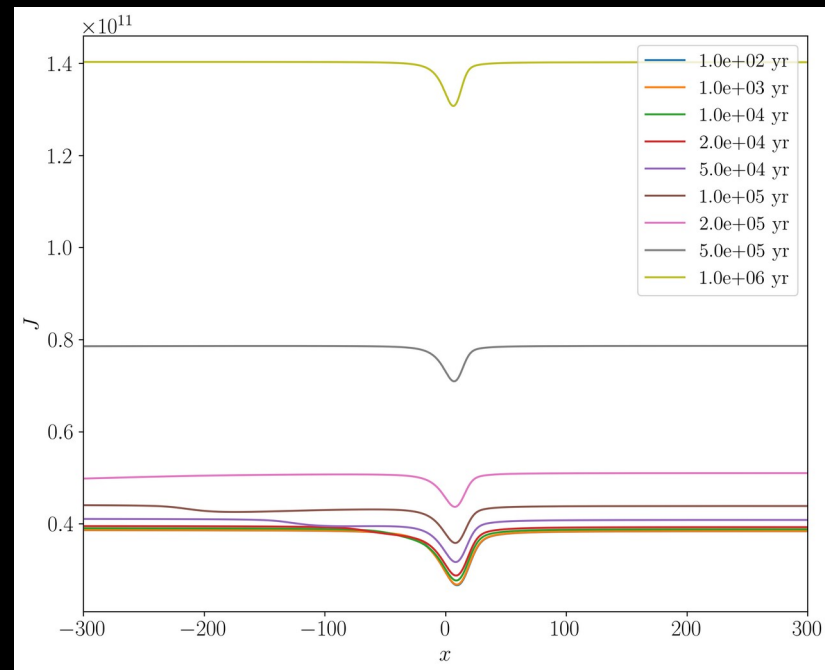
Fiducial Continuum Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



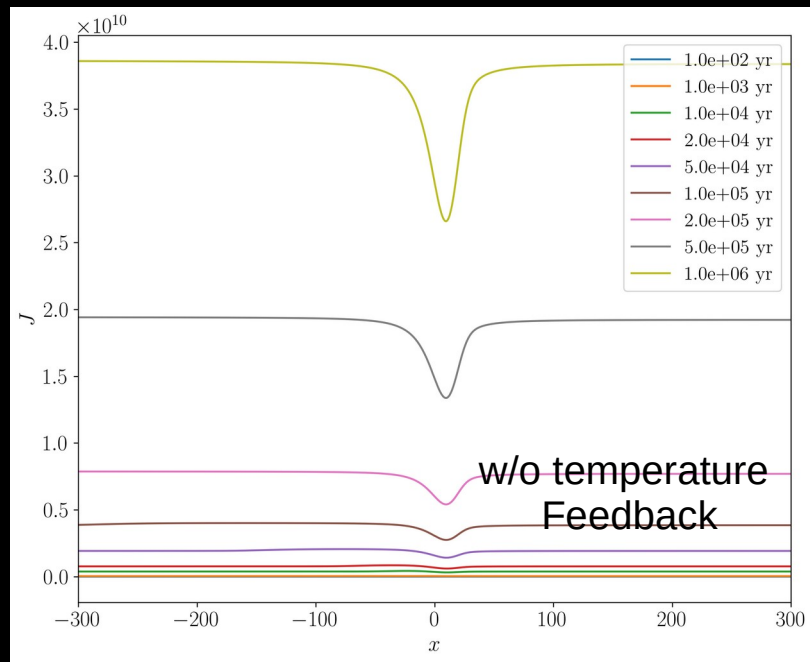
Fiducial Continuum Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



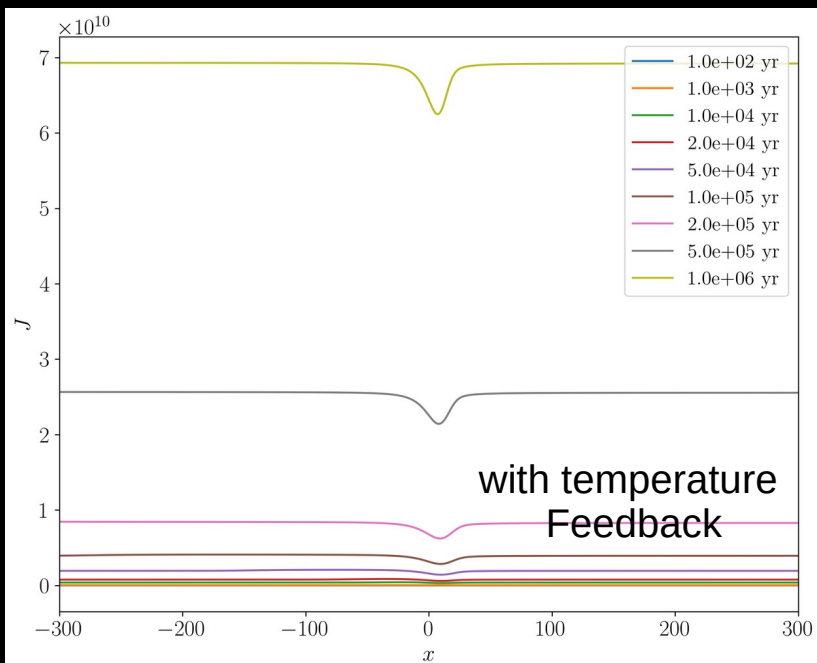
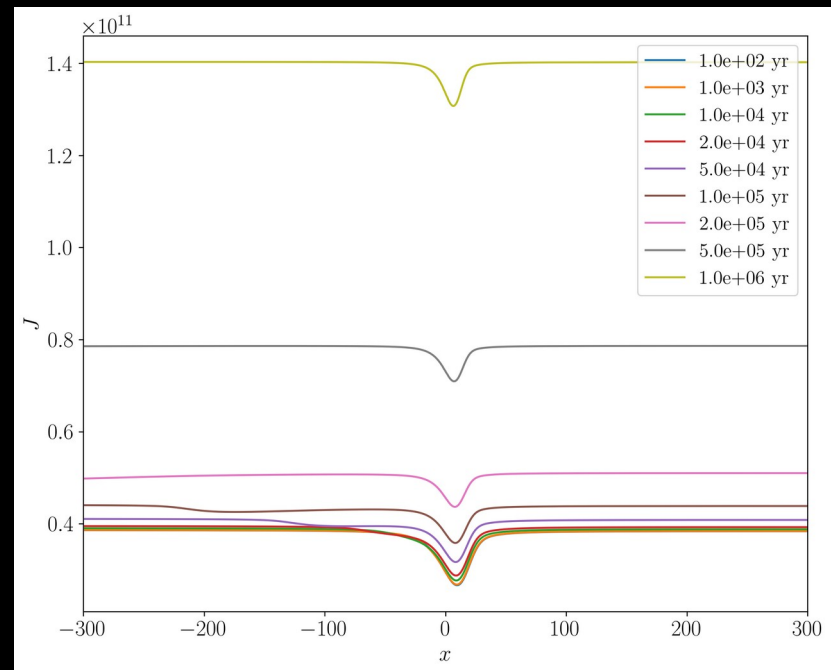
Temperature Feedback starts



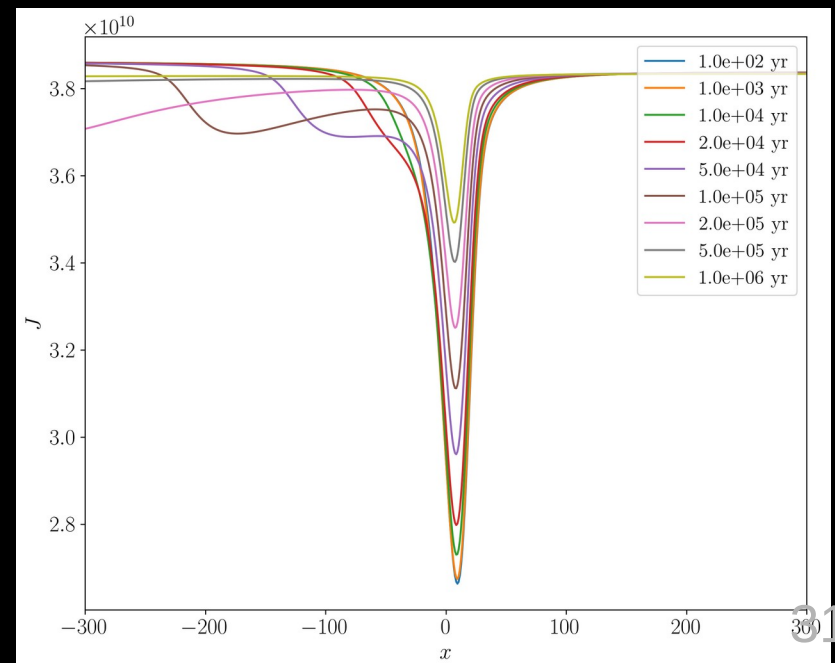
Fiducial Continuum Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



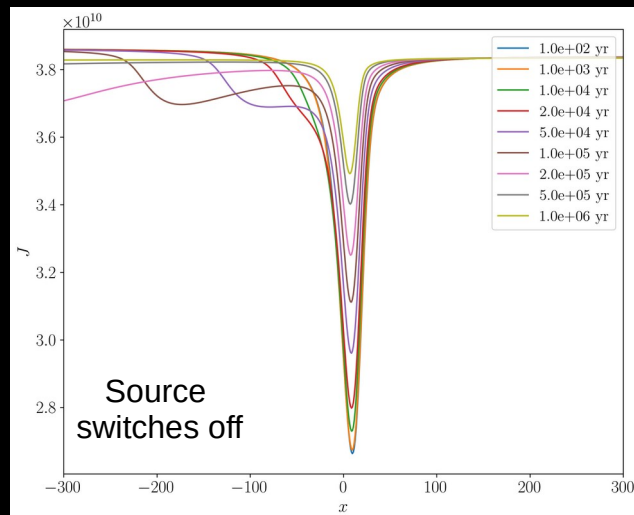
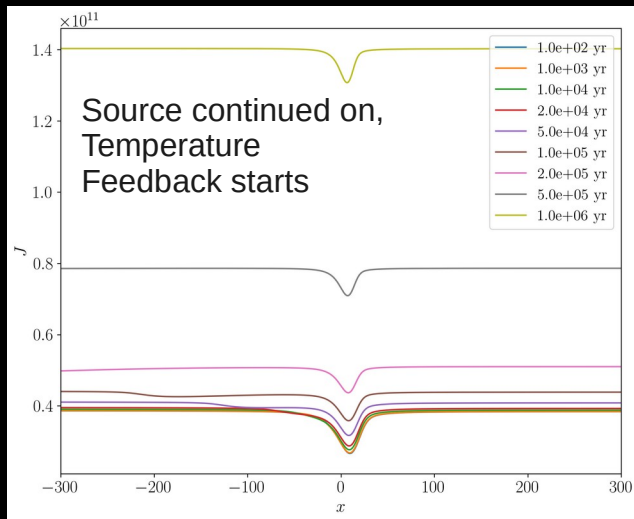
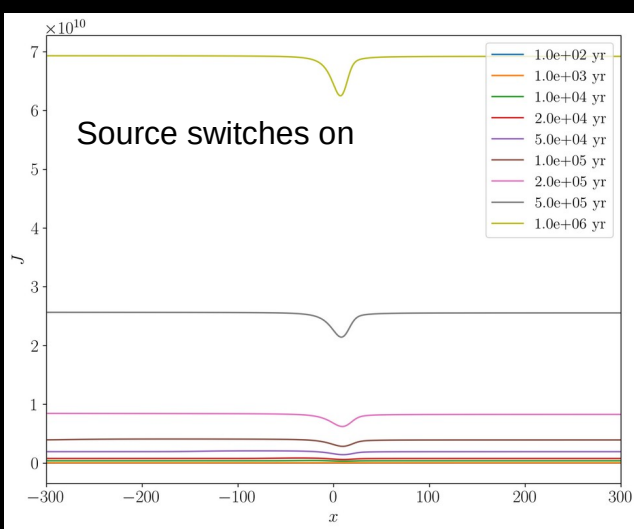
Temperature Feedback starts



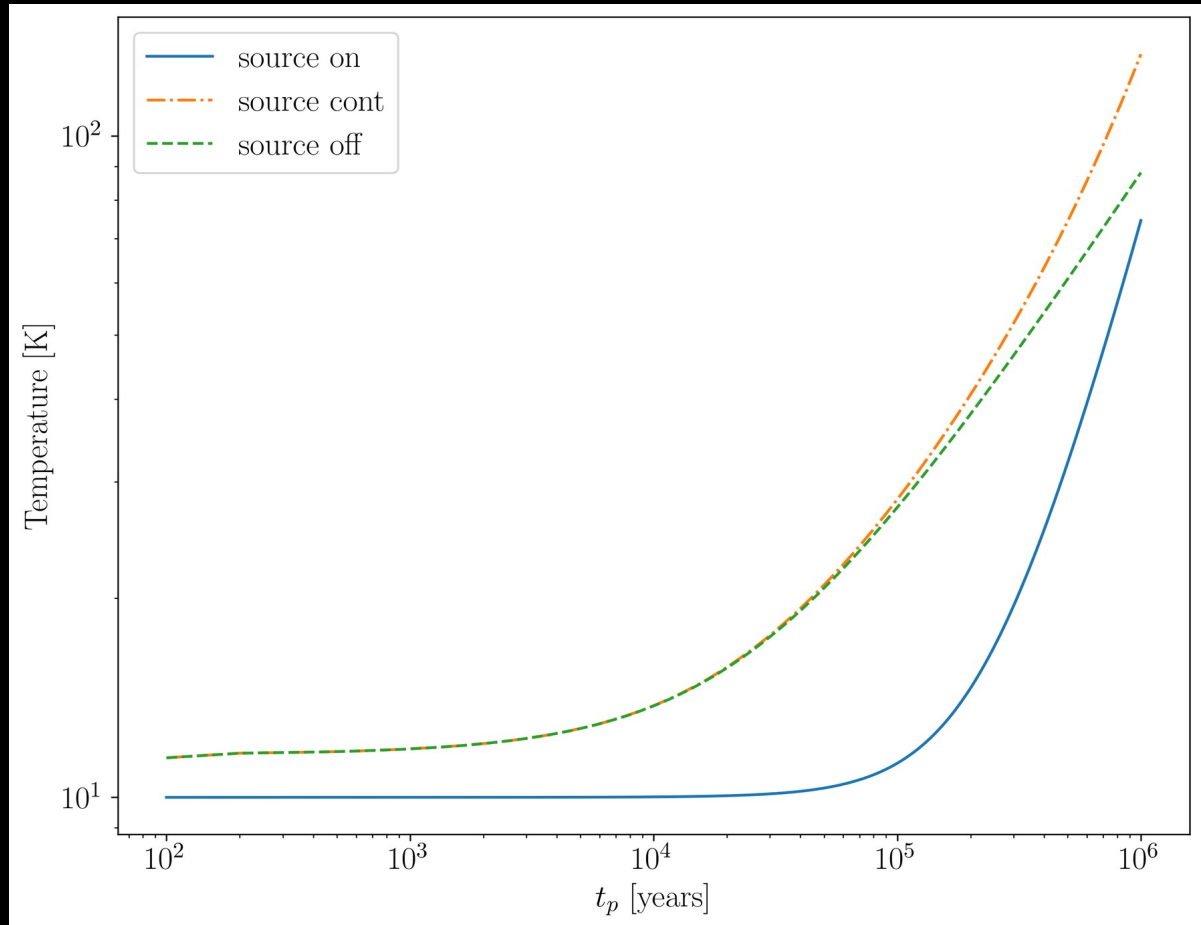
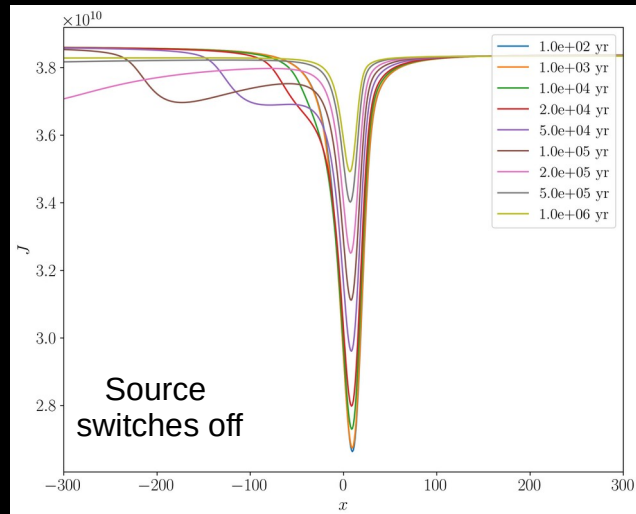
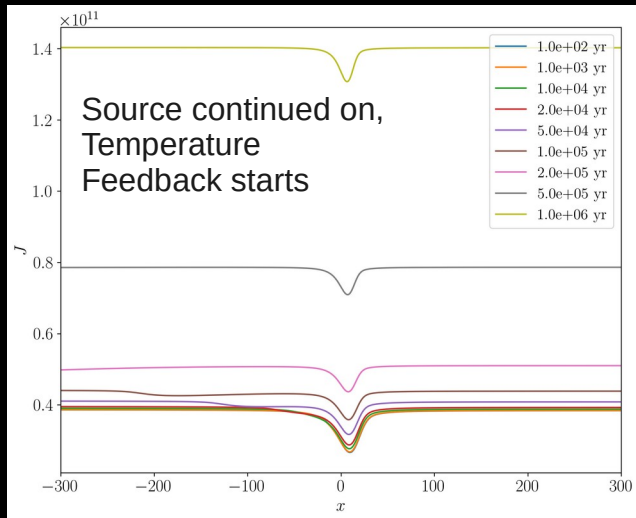
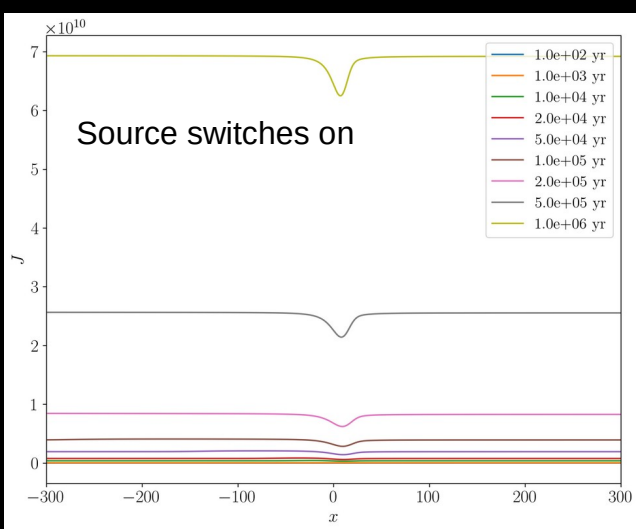
Source switches off



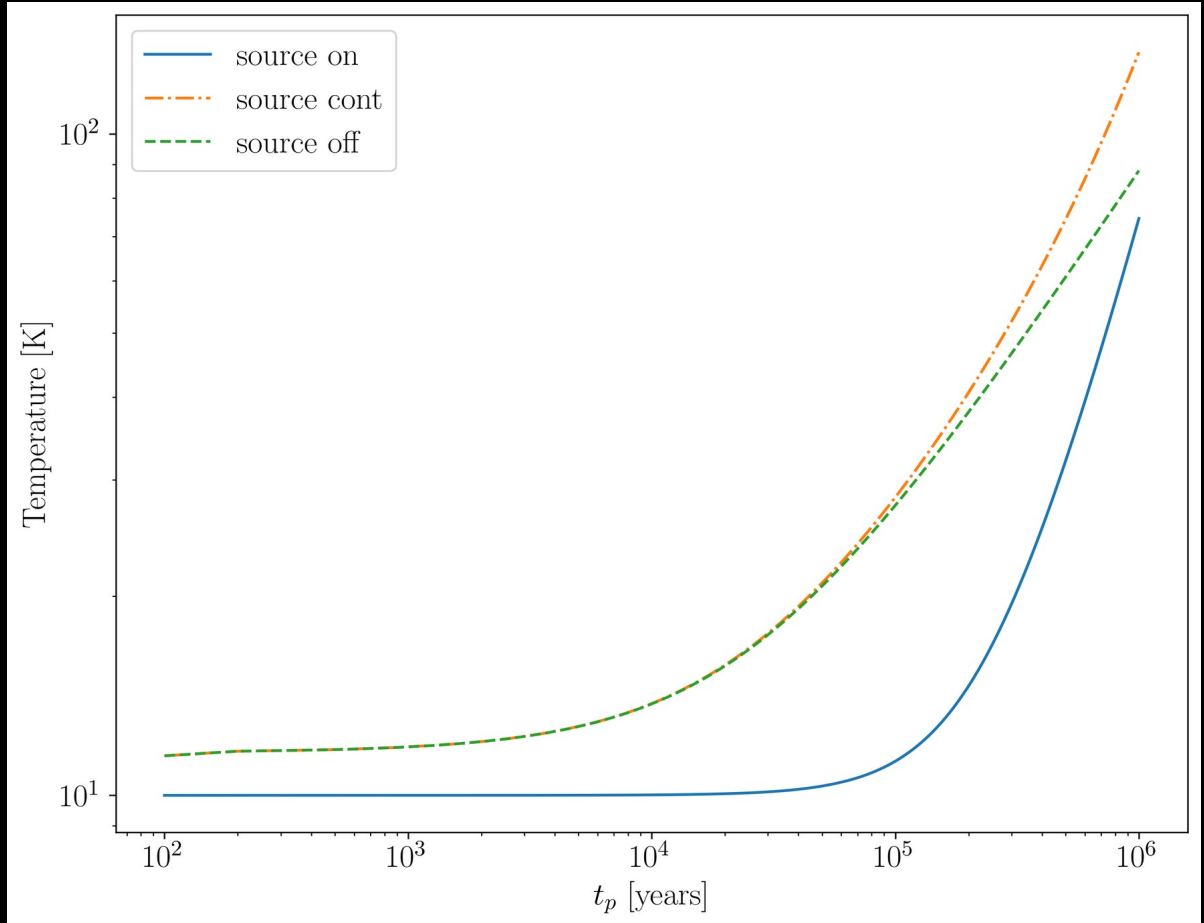
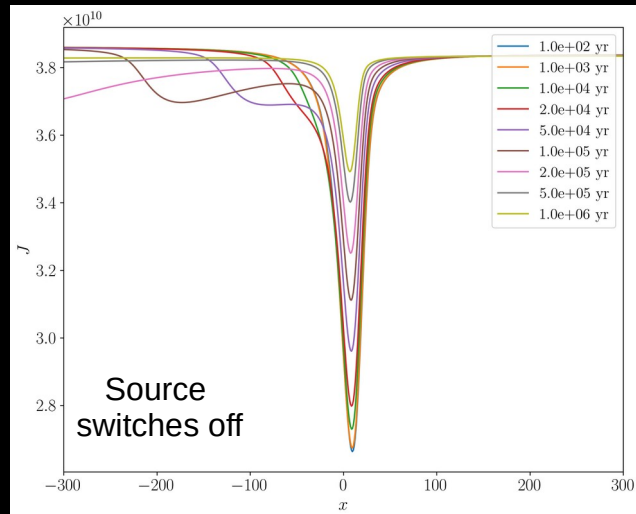
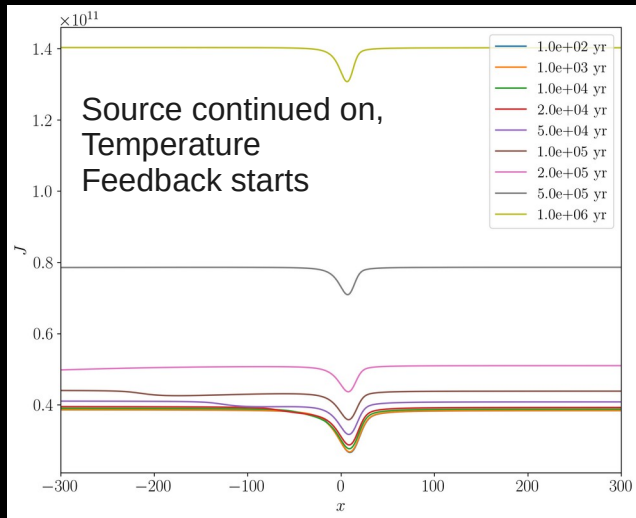
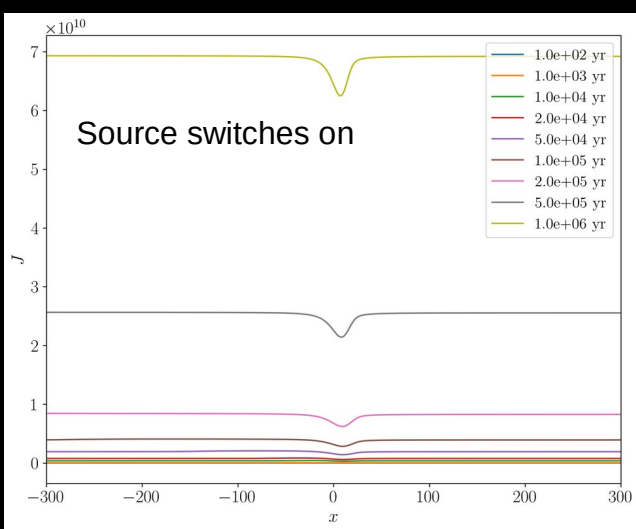
Fiducial Continuum Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



Fiducial Continuum Ly- α profile evolution for $z = 20$ and $T_K = 10$ K

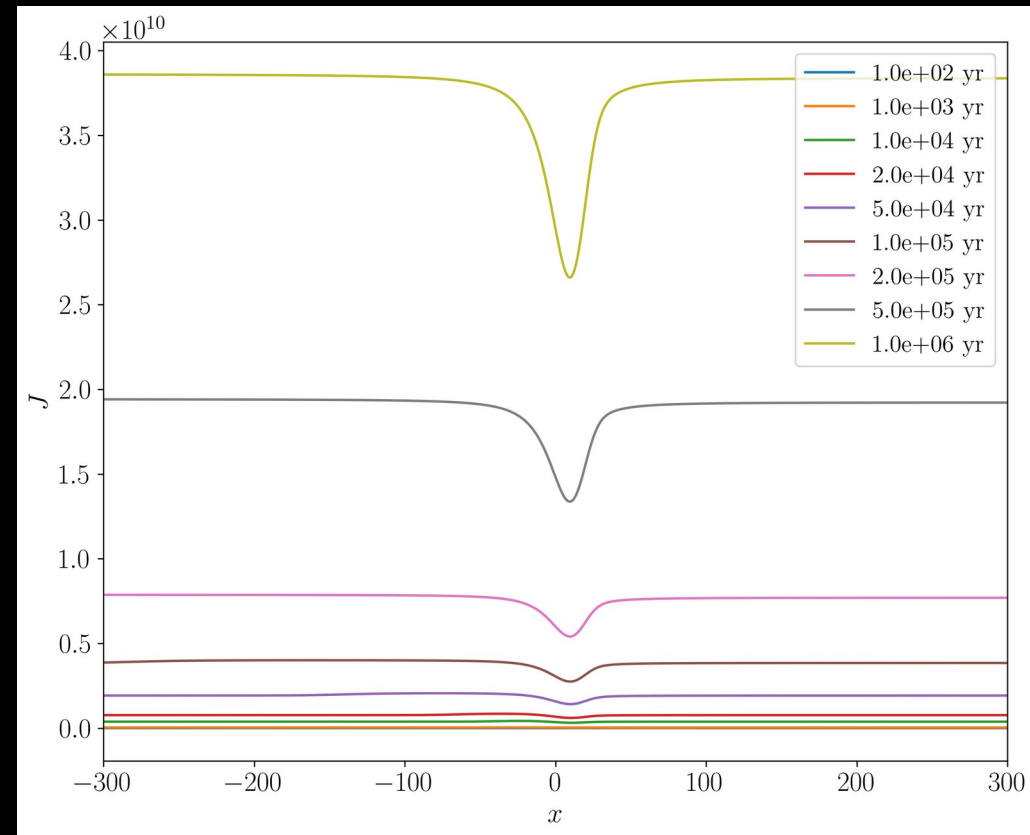
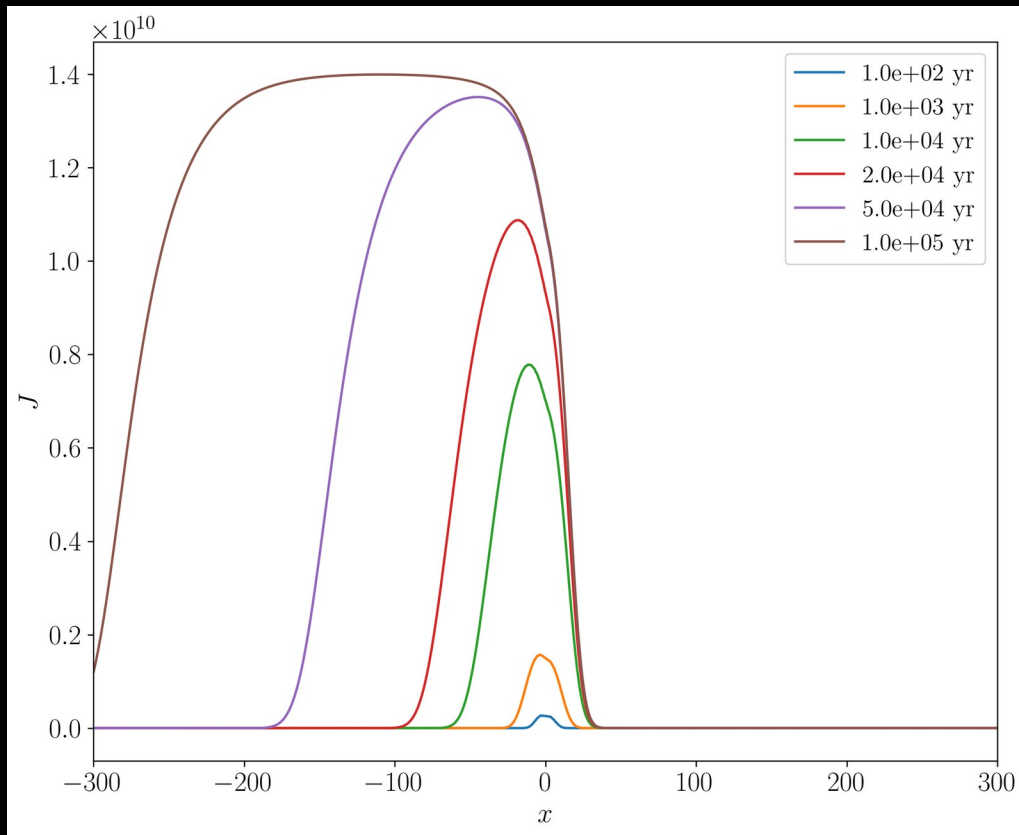


Fiducial Continuum Ly- α profile evolution for $z = 20$ and $T_K = 10$ K



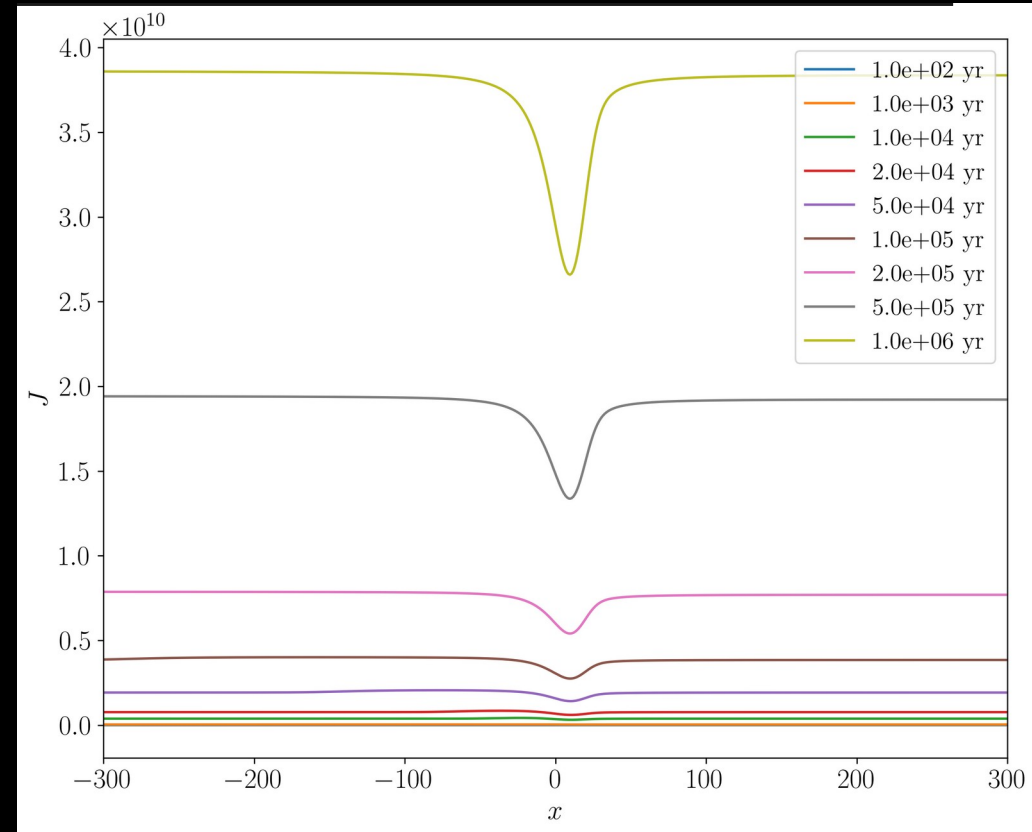
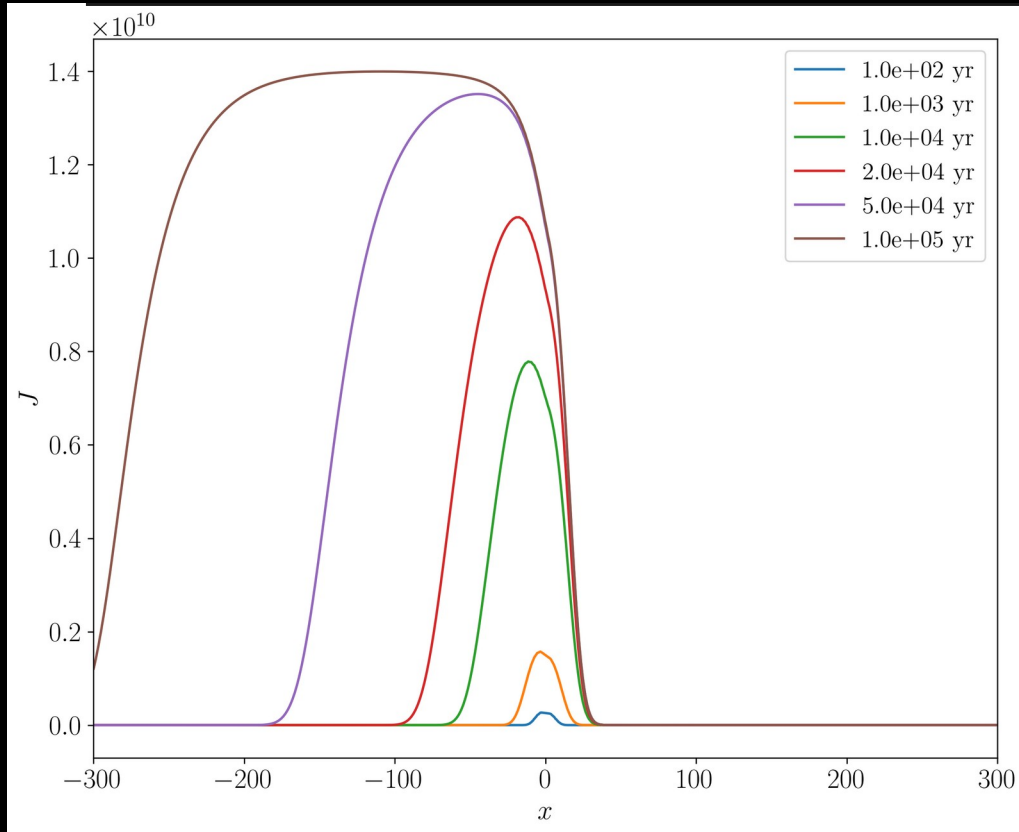
- Temperature feedback, source switching on, switching off delay/prevent quasi-static solution
- Continuum photons cause overall less heating than previously computed in literature.

When does Lyman- α photons profile reach quasi-static state?



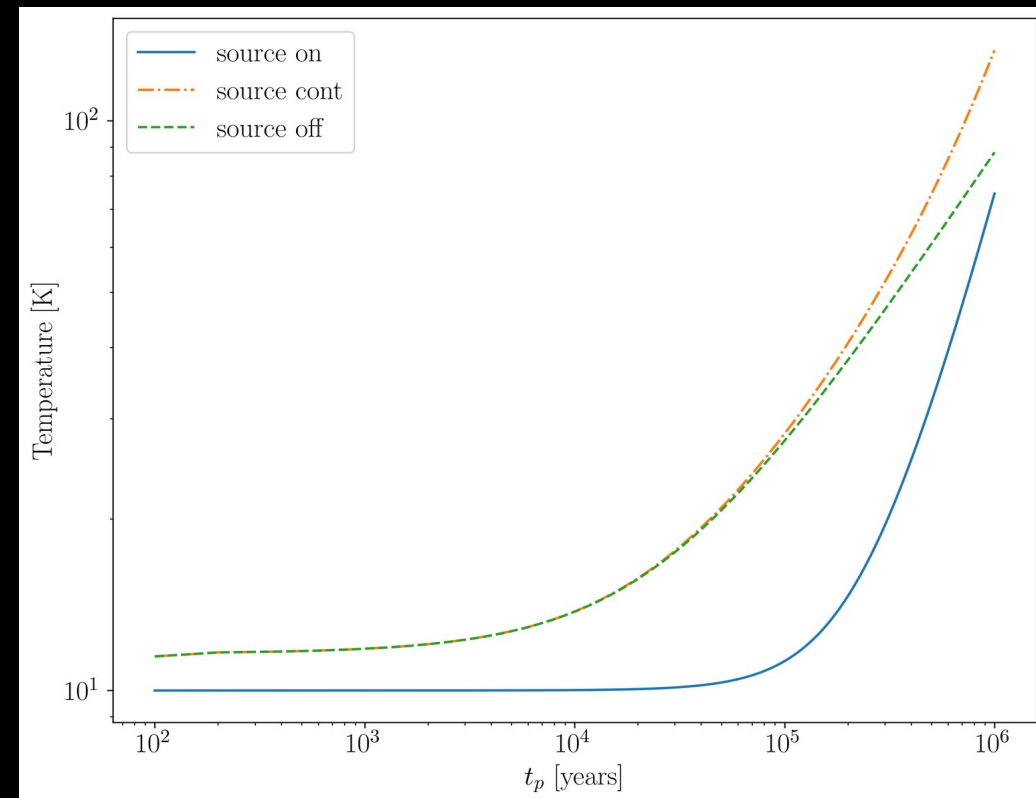
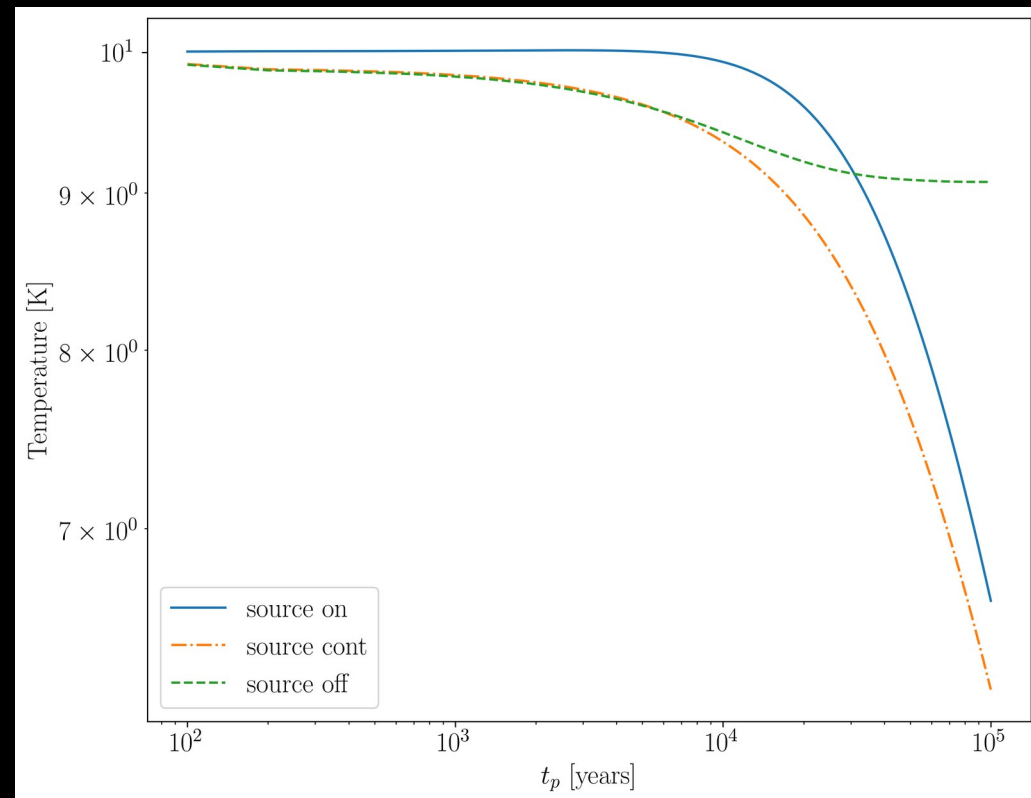
- Fiducial continuum Lyman- α photons profile evolution for $z = 20$ and $T_K = 10$ K

What is the lifetime of the Lyman- α photons Sources?



- Fiducial continuum Lyman- α photons profile evolution for $z = 20$ and $T_K = 10$ K
- Life span of Pop III stars $\sim 0.1-10$ Myr

How do the source and temperature feedback affect temperature evolution?



- Fiducial continuum Lyman- α photons profile evolution for $z = 20$ and $T_K = 10$ K
- Life span of Pop III stars ~ 0.1 - 10 Myr
- Continuum photons cause overall less heating than previously computed in literature. Injected photons cause less cooling.

Summary

- Lyman- α continuum photons heat up and injected photons cool the IGM during cosmic dawn – for quasi-static solution.
- Temperature feedback, source switching on and off delay or prevent quasi-static solution.
- The time scale to reach quasi-static solution is of the same order of magnitude as the source lifetime of PopIII stars during CD.
- Due to all these effect continuum photons cause overall less heating than previously computed in literature. Injected photons cause overall less cooling.
- **Future Work:** Combine effects of continuum and injected photons.

Thank you!